

# Singular Liouville Equations on $S^2$ : Sharp Inequalities and Existence Results

Gabriele Mancini\*

## Abstract

We prove a sharp Onofri-type inequality and non-existence of extremals for a Moser-Trudinger functional on  $S^2$  in the presence of potential having positive order singularities. We also investigate the existence of critical points and give some sufficient conditions under symmetry or nondegeneracy assumptions.

## 1 Introduction

In this work we study sharp Onofri-type inequalities on the standard Euclidean sphere  $(S^2, g_0)$ , and existence of critical points for a singular Moser-Trudinger functional. Given a smooth, closed surface  $\Sigma$ , and  $m$  points  $p_1, \dots, p_m \in \Sigma$ , we consider the functional

$$J_\rho^h(u) = \frac{1}{2} \int_\Sigma |\nabla u|^2 dv_g + \frac{\rho}{|\Sigma|} \int_\Sigma u dv_g - \rho \log \left( \frac{1}{|\Sigma|} \int_\Sigma h e^u dv_g \right) \quad (1)$$

where  $h$  is a positive singular potential satisfying

$$h \in C^\infty(\Sigma \setminus \{p_1, \dots, p_m\}) \quad \text{and} \quad h(x) \approx d(x, p_i)^{2\alpha} \text{ with } \alpha_i > -1 \text{ near } p_i, \quad (2)$$

$i = 1, \dots, m$ . Functionals of this kind were first introduced, for the regular case  $m = 0$ , by Moser ([18], [19]), in connection to the study of the Gaussian curvature equation on compact surfaces and Nirenberg's problem on  $S^2$ . They also have a role in spectral analysis due to Polyakov's formula (see [23], [24], [22], [21]). In the case  $m > 0$ , the functional (1) appears in the problem of prescribing the Gaussian curvature of Riemannian metrics with conical singularities. We recall that a metric on  $\Sigma$  with conical singularities of order  $\alpha_1, \dots, \alpha_m > -1$  in  $p_1, \dots, p_m$ , is a metric of the form  $e^u g$  where  $g$  is smooth metric on  $\Sigma$ , and  $u \in C^\infty(\Sigma \setminus \{p_1, \dots, p_m\})$  satisfies

$$|u(x) + 2\alpha_i \log d(x, p_i)| \leq C \quad \text{near } p_i, \quad i = 1, \dots, m.$$

---

\*S.I.S.S.A./I.S.A.S, Via Bonomea 265, 34136 Trieste (Italy) - gmancini@sissa.it

The author is supported by the PRIN project *Variational and perturbative aspects of nonlinear differential problems*.

It is possible to prove (see Proposition 2.1 in [3]) that a metric of this form has Gaussian curvature  $K$  if and only if  $u$  is a distributional solution of the Gaussian curvature equation

$$-\Delta_g u = 2Ke^u - 2K_g - 4\pi \sum_{i=1}^m \alpha_i \delta_{p_i}. \quad (3)$$

where  $K_g$  is the Gaussian curvature of  $(\Sigma, g)$ . If  $\chi(\Sigma) + \sum_{i=1}^m \alpha_i \neq 0$  and  $K_g$  is constant, (3) is equivalent to the singular Liouville equation

$$-\Delta_g u = \rho \left( \frac{Ke^u}{\int_{\Sigma} Ke^u dv_g} - \frac{1}{|\Sigma|} \right) - 4\pi \sum_{i=1}^m \alpha_i \left( \delta_{p_i} - \frac{1}{|\Sigma|} \right) \quad (4)$$

for

$$\rho = \rho_{geom} := 4\pi \left( \chi(\Sigma) + \sum_{i=1}^m \alpha_i \right). \quad (5)$$

Denoting by  $G$  the Green's function of  $(\Sigma, g)$ , that is the solution of

$$\begin{cases} -\Delta_g G(x, \cdot) = \delta_x & \text{on } \Sigma \\ \int_{\Sigma} G(x, y) dv_g(y) = 0, \end{cases} \quad (6)$$

the change of variable  $u \longleftrightarrow u + 4\pi \sum_{i=1}^m \alpha_i G(x, p_i)$  reduces (4) to

$$-\Delta_g u = \rho \left( \frac{he^u}{\int_{\Sigma} he^u dv_g} - \frac{1}{|\Sigma|} \right) \quad (7)$$

that is the Euler-Lagrange equation of the functional (1) corresponding to the potential

$$h(x) = Ke^{-4\pi \sum_{i=1}^m \alpha_i G_{p_i}}, \quad (8)$$

which satisfies (2). Equations (4) and (7) have also been widely studied in mathematical physics. For example, they appear in the description of Abelian vortices in Chern-Simons-Higgs theory, and have applications in Superconductivity and Electroweak theory ([26], [14]). We refer to [4], [7], [8], [16], [6], [12], [13], for some recent existence results.

A fundamental role in the variational analysis of (1) is played by singular versions of the standard Moser-Trudinger inequality (see [18, 27]). In [27], Troyanov proved that, for every function  $h$  satisfying (2), there exists a constant  $C = C(h, g, \Sigma)$  such that

$$\log \left( \frac{1}{|\Sigma|} \int_{\Sigma} he^{u-\bar{u}} dv_g \right) \leq \frac{1}{16\pi(1+\bar{\alpha})} \int_{\Sigma} |\nabla u|^2 dv_g + C(h, \Sigma, g) \quad (9)$$

$\forall u \in H^1(\Sigma)$ , where  $\bar{\alpha} = \min \left\{ 0, \min_{1 \leq i \leq m} \alpha_i \right\}$ . In particular the functional  $J_{\rho}^h$  is bounded from below  $\forall \rho \in (0, 8\pi(1+\bar{\alpha})]$  and is coercive for  $\rho \in (0, 8\pi(1+\bar{\alpha}))$ . Furthermore, it is possible to prove that the constant  $\frac{1}{16\pi}$  is sharp, that is

$$\inf_{H^1(\Sigma)} J_{\rho}^h = -\infty \quad \forall \rho > 8\pi(1+\bar{\alpha}).$$

In the special case  $m = 0$ ,  $h \equiv 1$  and  $(\Sigma, g) = (S^2, g_0)$ , a sharp version of (9) was proved by Onofri in [20]:

**Theorem A (Onofri's inequality [20]).**  $\forall u \in H^1(S^2)$  we have

$$\log \left( \frac{1}{4\pi} \int_{S^2} e^{u-\bar{u}} dv_{g_0} \right) \leq \frac{1}{16\pi} \int_{S^2} |\nabla u|^2 dv_{g_0}$$

with equality holding if and only if  $e^u g_0$  is a metric on  $S^2$  with positive constant Gaussian curvature, or, equivalently,  $u = \log |\det d\varphi| + c$  with  $c \in \mathbb{R}$  and  $\varphi : S^2 \rightarrow S^2$  a conformal diffeomorphism of  $S^2$ .

Motivated by this result, in [17] we started the study Onofri-type inequalities and existence of energy-minimizing solutions on  $S^2$  for the potential

$$h(x) = e^{-4\pi \sum_{i=1}^m \alpha_i G(x, p_i)}$$

(i.e. (8) with  $K \equiv 1$ ), and we extended Theorem A to the cases  $m = 1$ , and  $m = 2$  with  $\min\{\alpha_1, \alpha_2\} < 0$ .

**Theorem B ([17]).** If  $h = e^{-4\pi\alpha G_p}$  with  $\alpha \neq 0$ , then  $\forall u \in H^1(\Sigma)$

$$\log \left( \frac{1}{4\pi} \int_{S^2} h e^{u-\bar{u}} dv_{g_0} \right) < \frac{1}{16\pi \min\{1, 1+\alpha\}} \int_{S^2} |\nabla u|^2 dv_{g_0} + \max\{\alpha, -\log(1+\alpha)\}.$$

Moreover equation (7) has no solution for  $\rho = 8\pi \min\{1, 1+\alpha\}$ .

**Theorem C ([17]).** If  $h = e^{-4\pi\alpha_1 G_p - 4\pi\alpha_2 G_{p_2}}$  with  $p_2 = -p_1$ ,  $\alpha_1 = \min\{\alpha_1, \alpha_2\} < 0$ , then  $\forall u \in H^1(\Sigma)$

$$\log \left( \frac{1}{4\pi} \int_{S^2} h e^{u-\bar{u}} dv_{g_0} \right) \leq \frac{1}{16\pi(1+\alpha_1)} \int_{S^2} |\nabla u|^2 dv_{g_0} + \alpha_2 - \log(1+\alpha_1).$$

If  $\alpha_1 \neq \alpha_2$  there is no function realizing equality and no solution of (7) for  $\rho = 8\pi(1+\alpha_1)$ , while if  $\alpha_1 = \alpha_2$  then equality holds for  $u$  if and only if the following equivalent conditions are satisfied:

- $u$  is a solution of (7) for  $\rho = 8\pi(1+\alpha_1)$ .
- $h e^u g$  is a metric with constant positive Gaussian curvature and conical singularities of order  $\alpha_i$  in  $p_i$ ,  $i = 1, 2$ .
- If  $\pi$  denotes the stereographic projection from  $p_1$ , then

$$u \circ \pi^{-1}(y) = 2 \log \left( \frac{(1+|y|^2)^{1+\alpha_1}}{1+e^\lambda |y|^{2(1+\alpha_1)}} \right) + c$$

for some  $\lambda, c \in \mathbb{R}$ .

We stress that the critical parameter  $\rho = 8\pi(1+\bar{\alpha})$  is generally different from the geometric parameter (5) (except for some special cases, for example  $m = 2$  and  $\alpha_1 = \alpha_2 < 0$ ), thus critical points cannot always be interpreted in terms of metrics with prescribed curvature.

In this paper we will assume (8) with  $\alpha_i \geq 0$  for  $1 \leq i \leq m$  and

$$K \in C_+^\infty(S^2) := \{f \in C^\infty(S^2) : f(x) > 0 \quad \forall x \in S^2\}.$$

Our first result is a further extension of Onofri's inequality.

**Theorem 1.1.** *Assume that  $h$  satisfies (8) with  $K \in C_+^\infty(S^2)$  and  $\alpha_1, \dots, \alpha_m \geq 0$ , then*

$$\inf_{H^1(S^2)} J_{8\pi}^h = -8\pi \log \max_{S^2} h.$$

Moreover  $J_\rho^h$  has no minimum point, unless  $\alpha_1 = \dots = \alpha_m = 0$  (or, equivalently,  $m = 0$ ) and  $K$  is constant.

Clearly, the sharp value of the constant  $C(h, S^2, g_0)$  is given by

$$C(h, S^2, g_0) = -\frac{1}{8\pi(1+\bar{\alpha})} \inf_{H^1(S^2)} J_{8\pi(1+\bar{\alpha})}^h,$$

thus Theorem 1.1 is equivalent to the following sharp inequality:

**Corollary 1.1.** *If  $h$  satisfies (8) with  $K \in C_+^\infty(S^2)$  and  $\alpha_1, \dots, \alpha_m \geq 0$ , then  $\forall u \in H^1(S^2)$  we have*

$$\log \left( \frac{1}{4\pi} \int_{S^2} h e^{u-\bar{u}} dv_{g_0} \right) \leq \frac{1}{16\pi} \int_{S^2} |\nabla u|^2 dv_{g_0} + \log \max_{S^2} h$$

with equality holding if and only if  $m = 0$ ,  $K$  is constant and  $u$  realizes equality in Theorem A.

Theorem 1.1 states that  $J_{8\pi}^h$  has no minimum point, but does not exclude the existence of different kinds on critical points. In contrast to Theorem C, if  $\alpha_i > 0$  for  $1 \leq i \leq m$ , we will show that in many cases it is possible to find saddle points of  $J_{8\pi}^h$ . A simple example is given by the case in which  $h$  is axially symmetric. In this case an improved Moser-Trudinger inequality allows to minimize  $J_{8\pi}^h$  in the class of axially symmetric functions and find a solution of (7).

**Theorem 1.2.** *Assume that  $h$  satisfies (8) with  $m = 2$ ,  $p_1 = -p_2$ ,  $\min\{\alpha_1, \alpha_2\} = \alpha_1 > 0$  and  $K \in C_+^\infty(S^2)$  axially symmetric with respect to the direction identified by  $p_1$  and  $p_2$ . Then the Liouville equation (7) has an axially symmetric solution  $\forall \rho \in (0, 8\pi(1+\alpha_1))$ .*

In the last part of the paper we prove further general existence results using the Leray-Schauder degree theory introduced in [15], [10], [11], [12] and [13]. Solutions of (7) on the space

$$H_0 := \left\{ u \in H^1(S^2) : \int_{S^2} u dv_{g_0} = 0 \right\}. \quad (10)$$

can be obtained as solutions of  $T_\rho(u) + u = 0$  where  $T_\rho : H_0 \rightarrow H_0$  is defined by

$$T_\rho(u) = \Delta_{g_0}^{-1} \left( \frac{h e^u}{\int_{S^2} h e^u dv_{g_0}} - \frac{1}{4\pi} \right). \quad (11)$$

In [13], Chen and Lin computed the Leray-Schauder degree

$$d_\rho = \text{deg}_{LS}(Id + T_\rho, 0, B_R(0)) \quad (12)$$

for

$$\rho \notin \Gamma(\alpha_1, \dots, \alpha_n) := \left\{ 8\pi k_0 + 8\pi \sum_{i=1}^m k_i(1+\alpha_i) : k_0 \in \mathbb{N}, k_i \in \{0, 1\}, \sum_{i=0}^m k_i > 0 \right\}. \quad (13)$$

If  $m \geq 2$ , one has  $d_\rho \neq 0$  for any  $\rho \in (0, 8\pi(1+\bar{\alpha})) \setminus 8\pi\mathbb{N}$ . While Theorem 1.1 implies blow-up as  $\rho \nearrow 8\pi$ , we can find solutions for  $\rho = 8\pi$  by taking  $\rho \searrow 8\pi$ , provided the Laplacian of  $K$  is not too large at the critical points of  $h$ .

**Theorem 1.3.** *If  $h$  satisfies (8) with  $K \in C_+^\infty(S^2)$ ,  $m \geq 2$ ,  $\alpha_1, \dots, \alpha_m > 0$  and*

$$\Delta_{g_0} \log K(x) < \sum_{i=1}^m \alpha_i \quad (14)$$

$\forall x \in \Sigma$  such that  $\nabla h(x) = 0$ , then equation (7) has a solution for  $\rho = 8\pi$ .

The same strategy can be used for  $\rho = 8k\pi$ , with  $k < 1 + \alpha_1$ .

**Theorem 1.4.** *If  $h$  satisfies (8) with  $K \in C_+^\infty(S^2)$ ,  $m \geq 2$ ,  $0 < \alpha_1 \leq \dots \leq \alpha_m$  and*

$$\Delta_{g_0} \log K(x) < \sum_{i=1}^m \alpha_i + 2(1 - k) \quad (15)$$

$\forall x \in S^2$ , then equation (7) has a solution for  $\rho = 8k\pi$ ,  $k < 1 + \alpha_1$ .

Note that Theorems 1.3 and 1.4 can be applied in the case  $K \equiv 1$ . If the sign condition (14) is not satisfied, then it is not possible to exclude blow-up of solutions as  $\rho \rightarrow 8\pi$ . However, as it is pointed out in the introduction of [11], under some non-degeneracy assumptions on  $h$ , the Leray Schauder degree  $d_{8\pi}$  is well defined and can be explicitly computed by taking into account the contributions of all the blowing-up families of solutions. In particular one can prove that  $d_{8\pi} \neq 0$  under one of the following conditions.

**Theorem 1.5.** *Let  $h$  be a Morse function on  $S^2 \setminus \{p_1, \dots, p_m\}$  satisfying (8) with  $K \in C_+^\infty(S^2)$ ,  $m \geq 0$ ,  $\alpha_1, \dots, \alpha_m > 0$  and assume  $\Delta_{g_0} \log h \neq 0$  at all the critical points of  $h$ . If  $h$  has  $r$  local maxima and  $s$  saddle points in which  $\Delta_{g_0} h < 0$ , then equation (7) has a solution for  $\rho = 8\pi$  provided  $r \neq s + 1$ .*

**Theorem 1.6.** *Let  $h$  be a Morse function on  $S^2 \setminus \{p_1, \dots, p_m\}$  satisfying (8) with  $K \in C_+^\infty(S^2)$ ,  $m \geq 0$ ,  $\alpha_1, \dots, \alpha_m > 0$  and assume  $\Delta_{g_0} \log h \neq 0$  at all the critical points of  $h$ . If  $h$  has  $r'$  local minima in  $S^2 \setminus \{p_1, \dots, p_m\}$  and  $s'$  saddle points in which  $\Delta_{g_0} h > 0$ , then equation (7) has a solution for  $\rho = 8\pi$  provided  $s' \neq r' + \bar{d}$ , where*

$$\bar{d} := d_{8\pi+\varepsilon} = \begin{cases} 2 & m \geq 2, \\ 0 & m = 1, \\ -1 & m = 0. \end{cases}$$

In the regular case  $m = 0$ , Theorem 1.5 was first proved by Chang and Yang in [9] using a min-max scheme. A different proof was later given by Struwe [25] through a geometric flow approach.

## 2 Proof of the Main Results

The proof of Theorem 1.1 is a rather simple consequence of Theorem A.

*Proof of Theorem 1.1.* Let us consider

$$J_{8\pi}^1(u) := \frac{1}{2} \int_{S^2} |\nabla u|^2 dv_{g_0} + 2 \int_{S^2} u dv_{g_0} - 8\pi \log \left( \frac{1}{4\pi} \int_{S^2} e^u dv_{g_0} \right).$$

By Theorem A we have  $J_{8\pi}^1(u) \geq 0 \forall u \in H^1(S^2)$ . The condition  $\alpha_1, \dots, \alpha_m \geq 0$  guarantees  $h \in C^0(S^2)$ . Thus we have

$$\begin{aligned} J_{8\pi}^h(u) &\geq \frac{1}{2} \int_{S^2} |\nabla u|^2 dv_{g_0} + 2 \int_{S^2} u dv_{g_0} - 8\pi \log \left( \frac{1}{4\pi} \max_{S^2} h \int_{S^2} e^u dv_{g_0} \right) = \\ &= J_{8\pi}^1(u) - 8\pi \log \max_{S^2} h \geq -8\pi \log \max_{S^2} h. \end{aligned} \quad (16)$$

Since  $e^u > 0$  on  $S^2$ , equality can hold only if

$$h \equiv \max_{S^2} h$$

which, by (8), is possible only if  $\alpha_1 = \dots = \alpha_m = 0$  and  $K$  is constant. To complete the proof it is sufficient to observe that the lower bound in (16) is sharp. Let us fix a point  $p \in S^2$  such that  $h(p) = \max_{S^2} h$ , and consider the stereographic projection  $\pi : S^2 \setminus \{p\} \rightarrow \mathbb{R}^2$ . For  $t > 0$  we define  $u_t := \log |\det d\varphi_t|$ , where  $\{\varphi_t\}_{t>0}$  is the family of conformal diffeomorphisms of  $S^2$  that, in the local coordinates determined by  $\pi$ , corresponds to the family of dilations of  $\mathbb{R}^2$ , namely

$$\pi(\varphi_t(\pi^{-1}(y))) = ty \quad \forall y \in \mathbb{R}^2.$$

By Theorem A, we have  $J_{8\pi}^1(u_t) = 0 \forall t > 0$ . Moreover it is straightforward to verify that

$$\int_{S^2} e^{u_t} dv_{g_0} = 4\pi$$

and  $e^{u_t} \rightarrow 4\pi\delta_p$  weakly as measures on  $S^2$  for  $t \rightarrow \infty$ . Thus, one has

$$J_{8\pi}^h(u_t) = -8\pi \log \left( \frac{1}{4\pi} \int_{S^2} h e^{u_t} dv_{g_0} \right) \xrightarrow{t \rightarrow \infty} -8\pi \log h(p) = -8\pi \log \max_{S^2} h.$$

□

Let us now focus on the case of two antipodal singular points  $p_1 = -p_2$ . Given any point  $p \in S^2 \subset \mathbb{R}^3$  we consider the space

$$H_{rad,p} := \{u \in H^1(S^2) : \exists \varphi : [-1, 1] \rightarrow \mathbb{R} \text{ measurable s.t. } u(x) = v(x \cdot p) \text{ for a.e. } x \in S^2\}.$$

**Lemma 2.1.** *Suppose  $m = 2$ ,  $\min\{\alpha_1, \alpha_2\} = \alpha_1 > 0$  and  $p_2 = -p_1$ . If  $h$  is a positive function satisfying (2), then the Moser-Trudinger functional  $J_\rho^h$  is bounded from below on  $H_{rad,p_1}$  for any  $\rho \in (0, 8\pi(1 + \alpha_1))$ .*

*Proof.* Let us consider

$$\tilde{h}(x) := e^{-4\pi\alpha_1(G(x,p_1)+G(x,p_2))}.$$

Since  $h = Ke^{-4\pi\alpha_1G(x,p_1)-4\pi\alpha_2G(x,p_2)} \leq \tilde{h} \max_{x \in S^2} K(x)e^{4\pi(\alpha_1-\alpha_2)G(x,p_2)}$  it is sufficient to prove that the functional

$$\tilde{J}_\rho(u) := J_\rho^{\tilde{h}}(u) = \frac{1}{2} \int_{S^2} |\nabla u|^2 dv_{g_0} + \frac{\rho}{4\pi} \int_{S^2} u dv_{g_0} - \rho \log \left( \frac{1}{4\pi} \int_{S^2} \tilde{h} e^u dv_{g_0} \right)$$

is bounded from below for any  $\rho < 8\pi(1+\alpha_1)$ . Let us consider Euclidean coordinates  $(x_1, x_2, x_3)$  on  $S^2$  such that  $p_1 = (0, 0, -1)$ ,  $p_2 = (0, 0, 1)$ , and let  $\pi$  be the stereographic projection from the point  $p_2$ . Given a function  $u \in H^1(S^2)$  we define  $v(|y|) := (u(\pi^{-1}(y)))$ ,  $v_{\alpha_1}(y) := v(|y|^{\frac{1}{1+\alpha_1}})$  and  $u_{\alpha_1}(x) := v_{\alpha_1}(|\pi(x)|)$ . Then we have

$$\int_{S^2} |\nabla u|^2 dv_{g_0} = 2\pi \int_0^\infty t|v'(t)|^2 dt = (1+\alpha_1) \int_0^{+\infty} s|v'_{\alpha_1}(s)|^2 ds = (1+\alpha_1) \int_{S^2} |\nabla u_{\alpha_1}|^2 dv_{g_0}, \quad (17)$$

and, using that  $\sup_{t>0} \frac{1+t^{2(1+\alpha_1)}}{(1+t^2)^{1+\alpha_1}} < +\infty$ ,

$$\begin{aligned} \int_{S^2} \tilde{h} e^u dv_{g_0} &= 8\pi \int_0^{+\infty} e^{2\alpha_1} \frac{t^{2\alpha_1+1} e^{v(t)}}{(1+t^2)^{2(1+\alpha_1)}} dt \leq c_{\alpha_1} \int_0^{+\infty} \frac{t^{2\alpha_1+1} e^{v_{\alpha_1}(t^{1+\alpha_1})}}{(1+t^{2(1+\alpha_1)})^2} dt = \\ &= 4\tilde{c}_{\alpha_1} \int_0^{+\infty} \frac{s e^{v_{\alpha_1}(s)}}{(1+s^2)^2} ds = \tilde{c}_{\alpha_1} \int_{S^2} e^{v_{\alpha_1}} dv_{g_0}. \end{aligned} \quad (18)$$

Finally,  $\forall \varepsilon > 0, t \in \mathbb{R}^+$

$$\begin{aligned} |v(t) - v_{\alpha_1}(t)| &\leq \left| \int_t^{t^{\frac{1}{1+\alpha_1}}} |v'(s)| ds \right| \leq \left| \int_t^{t^{\frac{1}{1+\alpha_1}}} s|v'(s)|^2 ds \right|^{\frac{1}{2}} \left| \frac{\alpha_1}{1+\alpha_1} \log t \right| \leq \\ &\leq \frac{\varepsilon}{4\pi} \|\nabla u\|_2^2 + c_{\varepsilon, \alpha_1} |\log t| \end{aligned}$$

from which

$$\left| \int_{S^2} u dv_{g_0} - \int_{\Sigma} u_{\alpha_1} dv_{g_0} \right| \leq 8\pi \int_0^{+\infty} \frac{|v(t) - v_{\alpha_1}(t)|}{(1+t^2)^2} dt \leq \varepsilon \|\nabla u\|_2^2 + C_{\varepsilon, \alpha_1}. \quad (19)$$

(17), (18), (19) and the Moser-Trudinger inequality (9) imply

$$\begin{aligned} \tilde{J}_\rho(u) &\geq (1+\alpha_1) \left( \frac{1}{2} - \rho \varepsilon \right) \int_{S^2} |\nabla u_{\alpha_1}|^2 dv_{g_0} + \rho \int_{S^2} u_{\alpha_1} dv_{g_0} - \rho \log \left( \frac{1}{4\pi} \int_{S^2} e^{u_{\alpha_1}} dv_{g_0} \right) - C_{\varepsilon, \alpha_1, \rho} = \\ &= (1+\alpha_1) \left( \left( \frac{1}{2} - \rho \varepsilon \right) \int_{S^2} |\nabla u_{\alpha_1}|^2 dv_{g_0} - \frac{\rho}{1+\alpha_1} \log \left( \frac{1}{4\pi} \int_{S^2} e^{u_{\alpha_1} - \bar{u}_{\alpha_1}} dv_{g_0} \right) \right) - C_{\varepsilon, \alpha_1, \rho} \geq -\tilde{C}_{\varepsilon, \alpha_1, \rho} \end{aligned}$$

if  $\rho < 8\pi(1+\alpha_1)$  and  $\varepsilon$  is sufficiently small.  $\square$

**Remark 2.1.** *Arguing as in [17], it is possible to describe the behavior of sequences of minimum points of  $J_\rho^h$  in  $H_{rad,p_1}^1(S^2)$  as  $\rho \nearrow 8\pi(1+\alpha_1)$  to prove that also  $J_{8\pi(1+\alpha_1)}^h$  is bounded from below. Moreover if  $K \equiv 1$  and  $\alpha_1 = \alpha_2 = \alpha$  then we have*

$$\log\left(\frac{1}{4\pi}\int_{S^2} h e^{u-\bar{u}} dv_{g_0}\right) \leq \frac{1}{16\pi(1+\alpha)}\int_{S^2} |\nabla u|^2 dv_{g_0} + \alpha - \log(1+\alpha) \quad \forall u \in H_{rad,p_1}(S^2),$$

with equality holding for

$$u \circ \pi^{-1}(y) = 2 \log\left(\frac{(1+|y|^2)^{1+\alpha}}{1+e^\lambda|y|^{2(1+\alpha)}}\right) + c,$$

where  $\lambda, c \in \mathbb{R}$  and  $\pi$  is the stereographic projection from  $p_1$ .

*Proof of Theorem 1.2.* By Lemma 2.1,  $\forall \rho < 8\pi(1+\alpha_1) \exists \delta_\rho, C_\rho > 0$  such that

$$J_\rho^h(u) \geq \delta \int_{S^2} |\nabla u|^2 dv_{g_0} - C_\rho$$

$\forall u \in H_{rad,p_1}$ . Thus  $J_\rho^h$  is coercive on the space

$$\left\{ u \in H_{rad,p_1}, \int_\Sigma u dv_{g_0} = 0 \right\},$$

and by direct methods we can find a minimum point of  $J_\rho^h$  in  $H_{rad,p}^1$ . Since  $h \in H_{rad,p_1}^1$ , by Palais' criticality principle (see Remark 11.4 in [1]), this minimum point is a solution of (7).  $\square$

As a consequence of Theorems 1.1 and 1.2 we obtain a multiplicity result for equation (7). Indeed we can observe that if  $\rho < 8\pi$  is sufficiently close to  $8\pi$ , one has

$$\min_{u \in H^1(S^2)} J_\rho^h < \min_{u \in H_{rad,p_1}} J_\rho^h.$$

**Corollary 2.1.** *Suppose  $h$  satisfies the hypotheses of Theorem 1.2. There exists  $\varepsilon_0 > 0$  such that  $\forall \rho \in (8\pi - \varepsilon_0, 8\pi)$ , equation (7) has at least two solutions  $u, v$  such that  $u \in H_{rad,p_1}$  and  $v \in H^1(S^2) \setminus H_{rad,p_1}$ .*

*Proof.* For any  $\rho < 8\pi$  let us take two functions  $u_\rho \in H^1(S^2), v_\rho \in H_{rad,p_1}$ , such that

$$J_\rho^h(u_\rho) = \min_{H^1(S^2)} J_\rho^h, \quad J_\rho^h(v_\rho) = \min_{H_{rad,p_1}(S^2)} J_\rho^h(u) \quad \text{and} \quad \int_\Sigma u_\rho dv_{g_0} = \int_\Sigma v_\rho dv_{g_0} = 0.$$

We claim that, for  $\varepsilon$  sufficiently small and  $\rho \in (8\pi - \varepsilon, 8\pi)$ ,  $u_\rho \notin H_{rad,p_1}$  and in particular  $u_\rho \neq v_\rho$ . Assume by contradiction that there exists a sequence  $\rho_n \nearrow 8\pi$  for which  $u_{\rho_n} \in H_{rad,p_1}$ . Then, applying Lemma 2.1 as in the proof Theorem 1.2, we would have

$$J_{\rho_n}^h(u_{\rho_n}) \geq \delta \int_{S^2} |\nabla u_{\rho_n}|^2 dv_{g_0} - C$$

for some  $\delta, C > 0$ . Therefore  $\|\nabla u_{\rho_n}\|_2$  would be uniformly bounded and, up to subsequences,  $u_{\rho_n} \rightharpoonup u$  in  $H^1(S^2)$  with  $J_{8\pi}^h(u) = \inf_{H^1(S^2)} J_{8\pi}^h$ . This is not possible because we know by Theorem 1.1 that  $J_{8\pi}^h$  has no minimum point.  $\square$



Now we will discuss some sufficient conditions for the existence of solutions of (7), without symmetry assumptions on  $h$ . Let  $H_0$ ,  $T_\rho$ ,  $d_\rho$  and  $\Gamma(\alpha_1, \dots, \alpha_m)$  be defined as in (10), (11), (12) and (13). First of all we recall a well known result concerning blow-up analysis for sequences of solutions.

**Proposition 2.1** (See [2], [5]). *Let  $(\Sigma, g)$  be a compact Riemannian surface and let  $h$  be a function satisfying (8) with  $K \in C_+^\infty(\Sigma)$ . If  $u_n$  is a sequence of solutions of (7) on  $\Sigma$  with  $\rho = \rho_n \rightarrow \bar{\rho}$  and  $\int_\Sigma u_n dv_g = 0$ , Then, up to subsequences, one of the following holds:*

- (i)  $|u_n| \leq C$  with  $C$  depending only on  $\alpha_1, \dots, \alpha_m$ ,  $\max_\Sigma K$ ,  $\min_\Sigma K$  and  $\bar{\rho}$ .
- (ii) (blow-up). *There exists a finite set  $S = \{q_1, \dots, q_k\}$  such that  $u_n \rightarrow -\infty$  uniformly on compact subsets of  $\Sigma \setminus S$ . Moreover  $\frac{hu_n}{\int_\Sigma h e^{u_n} dv_g} \rightharpoonup \sum_{i=1}^k \beta_i \delta_{q_i}$  with  $\beta_i = 8\pi$  if  $q_i \in \Sigma \setminus \{p_1, \dots, p_m\}$  and  $\beta_i = 8\pi(1 + \alpha_j)$  if  $q_i = p_j$  for some  $1 \leq j \leq m$ .*

Clearly case (ii) is possible only if  $\bar{\rho} \in \Gamma(\alpha_1, \dots, \alpha_m)$ . As a direct consequence of Proposition 2.1 we get that, if  $E$  is a compact subset of  $(0, +\infty) \setminus \Gamma(\alpha_1, \dots, \alpha_m)$ , the set of all the solutions of (7) in  $H_0$  with  $\rho \in E$  is a bounded subset  $H_0$ . This bound depends only on  $E$ ,  $\alpha_1, \dots, \alpha_m$  and on  $\max_\Sigma K$ ,  $\min_\Sigma K$ , thus, using the homotopy invariance of the Leray-Schauder degree, one can prove that, if  $R$  is chosen sufficiently large,  $d_\rho$  is well defined and does not depend on  $R$  and  $K$ . Moreover  $d_\rho$  is constant on every connected component of  $(0, +\infty) \setminus \Gamma(\alpha_1, \dots, \alpha_m)$ . In [13] Chen and Lin introduced the generating function

$$g(x) := (1 + x + x^2 + x^3 \dots)^{m-2} \prod_{i=1}^m (1 - x^{1+\alpha_i})$$

and observed that

$$g(x) = 1 + \sum_{j=1}^{\infty} b_j x^{n_j} \tag{20}$$

where  $n_1 < n_2 < n_3 < \dots$  are such that

$$\Gamma(\alpha_1, \dots, \alpha_m) = \{8\pi n_j : j \geq 1\}.$$

**Theorem D** ([13]). *Let  $h$  be a function satisfying (8), then for  $\rho \in (8\pi n_k, 8\pi n_{k+1})$  we have*

$$d_\rho = \sum_{j=0}^k b_j$$

where  $b_0 = 1$  and  $b_j$  are the coefficients in (20).

As a consequence of this formula, (7) has a solution for any  $\rho \in (0, 8\pi(1 + \alpha_1)) \setminus 8\pi\mathbb{N}$ .

**Lemma 2.2.** *Suppose that  $h$  satisfies (8) with  $K \in C_+^\infty(S^2)$ ,  $m \geq 2$  and  $0 < \alpha_1 \leq \dots \leq \alpha_m$ . Then equation (7) has a solution  $\forall \rho \in (0, 8\pi(1 + \alpha_1)) \setminus 8\pi\mathbb{N}$ .*

*Proof.* Indeed the first negative coefficient appearing in the expansion

$$g(x) = (1 + x + x^2 + x^3 \dots)^{m-2} \prod_{i=1}^m (1 - x^{1+\alpha_i}) = 1 + \sum_{j=1}^{\infty} b_j x^{n_j}$$

is the coefficient of  $x^{1+\alpha_1}$ , i.e.

$$g(x) = \sum_{j=0}^{\infty} b_j x^{n_j}$$

with  $b_0 = 1$  and  $b_j \geq 0$  for any  $j \geq 1$  such that  $n_j < 1 + \alpha_1$ . From Theorem D it follows that  $d_\rho \geq 1$  for  $\rho \in (0, 8\pi(1 + \alpha_1)) \setminus 8\pi\mathbb{N}$ .  $\square$

**Remark 2.2.** *Lemma 2.2 only holds for  $m \geq 2$ . Indeed for  $m = 1$  and  $K \equiv 1$  one can use a Pohozaev-type identity (see [17]) to prove that (7) has no solutions for  $\rho \in [8\pi, 8\pi(1 + \alpha_1)]$ .*

**Remark 2.3.** *A different proof of Lemma 2.2 was given in [4] by Bartolucci and Malchiodi using topological methods.*

By Proposition 2.1, if  $\rho_n \rightarrow 8k\pi$  with  $k < 1 + \alpha_1$ , then any blowing-up sequence of solutions of (7) must concentrate around exactly  $k$  points  $q_1, \dots, q_k \in \Sigma \setminus \{p_1, \dots, p_m\}$ . A more precise description of the blow-up set is given in [10] (see also [12], [13]):

**Proposition 2.2** ([10], [12]). *Let  $u_n$  be a sequence of solutions of (7) with  $\rho = \rho_n \rightarrow 8\pi k$  and  $k < 1 + \alpha_1$ . If alternative (ii) of Proposition 2.1 holds, then  $u_n$  has exactly  $k$  blow-up points  $q_1, \dots, q_k \in \Sigma \setminus \{p_1, \dots, p_m\}$  and  $(q_1, \dots, q_k)$  is a critical point of the function*

$$f_h(x_1, \dots, x_k) := \sum_{j=1}^k \left( \log h(x_j) + \sum_{l \neq j} G(x_l, x_j) \right)$$

in

$$\{(x_1, \dots, x_k) \in (S^2)^k : x_i \neq x_j \text{ for } i \neq j\}.$$

Moreover we have

$$\rho_n - 8k\pi = \sum_{j=1}^k h(q_{j,n})^{-1} (\Delta_{g_0} \log h(q_{j,n}) + 2(k-1)) \frac{\lambda_{j,n}}{e^{\lambda_{j,n}}} + O(e^{-\lambda_{j,n}})$$

where  $q_{j,n}$  are the local maxima of  $u_n$  near  $q_j$  and  $\lambda_{j,n} = u_n(q_{j,n})$ .

*Proof of Theorems 1.3 and 1.4.* Take a sequence  $\rho_n \searrow 8k\pi$  and a solution  $u_n \in H_0$  of (7) for  $\rho = \rho_n$ . By Propositions 2.1, 2.2 and standard elliptic estimates, either  $u_n$  is uniformly bounded in  $W^{2,q}(S^2)$  for any  $q \geq 1$  or  $u_n$  blows-up at  $(q_1, \dots, q_k) \in \Sigma \setminus \{p_1, \dots, p_m\}$ . In the former case we have  $u_n \rightarrow u$  in  $H^1(S^2)$  and  $u$  satisfies (7) with  $\rho = 8\pi k$ . The latter case can be excluded using (14), (15). Indeed we have

$$\Delta_{g_0} \log h(q_j) + 2(k-1) = \Delta_{g_0} \log K - \sum_{i=1}^m \alpha_i + 2(k-1) < 0$$

for any  $j$ . Denoting  $q_{j,n}$  the maximum point of  $u_n$  near  $q_j$  and  $\lambda_{j,n} = u_n(q_{j,n})$ , by Proposition 2.2 we get

$$\begin{aligned} \rho_n - 8\pi k &= \sum_{j=1}^k h(q_{j,n})^{-1} (\Delta_{g_0} \log h(q_{j,n}) + 2(k-1)) \frac{\lambda_{j,n}}{e^{\lambda_{j,n}}} + O(e^{-\lambda_{j,n}}) = \\ &= \sum_{j=1}^k h(q_j)^{-1} (\Delta_{g_0} \log h(q_j) + 2(k-1)) \lambda_{j,n} e^{-\lambda_{j,n}} + o(\lambda_{j,n} e^{-\lambda_{j,n}}) < 0 \end{aligned}$$

which contradicts  $\rho_n \searrow 8k\pi$ .  $\square$

In order to prove Theorems 1.5, 1.6 we need to compute the Leray-Schauder degree for  $\rho = 8\pi$ .

**Lemma 2.3.** *Let  $h$  be a function satisfying (8) with  $K \in C_+^\infty(\Sigma)$  and  $\alpha_1, \dots, \alpha_m > 0$ . If  $\Delta_{g_0} h(q) \neq 0$  for any  $q \in \Sigma \setminus \{p_1, \dots, p_m\}$  critical point of  $h$ , then  $d_{8\pi}$  is well defined.*

*Proof.* It is sufficient to prove that the set of solutions of (7) in  $H_0$  with  $\rho = 8\pi$  is a bounded subset of  $H_0$ . Assume by contradiction that there exists  $u_n \in H_0$  solution of (7) for  $\rho = 8\pi$  such that  $\|u_n\|_{H_0} \rightarrow +\infty$ . By Propositions 2.1 and 2.2, there exists  $q \in \Sigma \setminus \{p_1, \dots, p_m\}$  such that  $u_n \rightarrow 8\pi\delta_q$ ,  $\nabla h(q) = 0$  and

$$0 = h(q_n)^{-1} \Delta_{g_0} \log h(q_n) \lambda_n e^{-\lambda_n} + O(e^{-\lambda_n}) = h(q)^{-2} \Delta_{g_0} h(q) \lambda_n e^{-\lambda_n} + o(\lambda_n e^{-\lambda_n})$$

where  $\lambda_n := \max_\Sigma u_n$  and  $u_n(q_n) = \lambda_n$ . Since  $\Delta_{g_0} h(q) \neq 0$  this is not possible.  $\square$

Under nondegeneracy assumptions, Chen and Lin proved that for any critical  $q$  point of  $h$  there exists a blowing-up sequence of solutions which concentrates at  $q$ . Moreover they were able to compute the total contribution to the Leray-Schauder degree of all the solutions concentrating at  $q$ .

**Proposition 2.3** (see [11], [13]). *Assume that  $h$  is a Morse function on  $\Sigma \setminus \{p_1, \dots, p_m\}$ . Given a critical point  $q \in \Sigma \setminus \{p_1, \dots, p_m\}$  of  $h$ , the total contribution to  $d_{8\pi}$  of all the solutions of (7) concentrating at  $q$  is equal to  $\text{sgn}(\rho - 8\pi)(-1)^{\text{ind}_p}$ , where  $\text{ind}_p$  is the Morse index of  $p$  as critical point of  $h$ .*

*Proof of Theorems 1.5, 1.6.* Let us denote

$$\Lambda_- = \{q \in \Sigma \setminus \{p_1, \dots, p_m\} : \nabla h(q) = 0, \Delta_{g_0} h(q) < 0\},$$

$$\Lambda_+ = \{q \in \Sigma \setminus \{p_1, \dots, p_m\} : \nabla h(q) = 0, \Delta_{g_0} h(q) > 0\}.$$

By Proposition 2.3 we have

$$d_{8\pi} = 1 - \sum_{q \in \Lambda_-} (-1)^{\text{ind}_q} = \bar{d} + \sum_{q \in \Lambda_+} (-1)^{\text{ind}_q},$$

where  $\bar{d}$  is the Leray-Schauder degree for  $\rho \in (8\pi, 8\pi + \varepsilon)$ . Clearly  $\Lambda_-$  contains only the local maxima of  $h$  and the saddle points of  $h$  in which  $\Delta_{g_0}h < 0$ , thus

$$d_{8\pi} = 1 - r + s.$$

Therefore we get existence of solutions if  $r \neq s + 1$ . Similarly we have

$$d_{8\pi} = \bar{d} - s' + r'$$

and we get solutions if  $s' \neq r' + \bar{d}$ .  $\bar{d}$  can be computed using Theorem D. If  $m \geq 2$ ,

$$g(x) = 1 + x + \dots \implies \bar{d} = 2.$$

If  $m = 1$  we have

$$g(x) := 1 - x - x^{1+\alpha} + x^{2(1+\alpha)} \implies \bar{d} = 0.$$

If  $m = 0$ , then

$$g(x) = 1 - 2x + x^2 \implies \bar{d} = -1.$$

This concludes the proof. □

## References

- [1] Antonio Ambrosetti and Andrea Malchiodi. *Nonlinear analysis and semilinear elliptic problems*, volume 104 of *Cambridge Studies in Advanced Mathematics*. Cambridge University Press, Cambridge, 2007.
- [2] D. Bartolucci and G. Tarantello. Liouville type equations with singular data and their applications to periodic multivortices for the electroweak theory. *Comm. Math. Phys.*, 229(1):3–47, 2002.
- [3] Daniele Bartolucci, Francesca De Marchis, and Andrea Malchiodi. Supercritical conformal metrics on surfaces with conical singularities. *Int. Math. Res. Not. IMRN*, (24):5625–5643, 2011.
- [4] Daniele Bartolucci and Andrea Malchiodi. An improved geometric inequality via vanishing moments, with applications to singular Liouville equations. *Comm. Math. Phys.*, 322(2):415–452, 2013.
- [5] Daniele Bartolucci and Eugenio Montefusco. Blow-up analysis, existence and qualitative properties of solutions for the two-dimensional Emden-Fowler equation with singular potential. *Math. Methods Appl. Sci.*, 30(18):2309–2327, 2007.
- [6] Alessandro Carlotto. On the solvability of singular Liouville equations on compact surfaces of arbitrary genus. *Trans. Amer. Math. Soc.*, 366(3):1237–1256, 2014.
- [7] Alessandro Carlotto and Andrea Malchiodi. A class of existence results for the singular Liouville equation. *C. R. Math. Acad. Sci. Paris*, 349(3-4):161–166, 2011.

- [8] Alessandro Carlotto and Andrea Malchiodi. Weighted barycentric sets and singular Liouville equations on compact surfaces. *J. Funct. Anal.*, 262(2):409–450, 2012.
- [9] Sun-Yung A. Chang and Paul C. Yang. Conformal deformation of metrics on  $S^2$ . *J. Differential Geom.*, 27(2):259–296, 1988.
- [10] Chiun-Chuan Chen and Chang-Shou Lin. Sharp estimates for solutions of multi-bubbles in compact Riemann surfaces. *Comm. Pure Appl. Math.*, 55(6):728–771, 2002.
- [11] Chiun-Chuan Chen and Chang-Shou Lin. Topological degree for a mean field equation on Riemann surfaces. *Comm. Pure Appl. Math.*, 56(12):1667–1727, 2003.
- [12] Chiun-Chuan Chen and Chang-Shou Lin. Mean field equations of Liouville type with singular data: sharper estimates. *Discrete Contin. Dyn. Syst.*, 28(3):1237–1272, 2010.
- [13] Chiun-Chuan Chen and Chang-Shou Lin. Mean field equation of Liouville type with singular data: topological degree. *Comm. Pure Appl. Math.*, 68(6):887–947, 2015.
- [14] Jooyoo Hong, Yoonbai Kim, and Pong Youl Pac. Multivortex solutions of the abelian Chern-Simons-Higgs theory. *Phys. Rev. Lett.*, 64(19):2230–2233, 1990.
- [15] Yan Yan Li. Harnack type inequality: the method of moving planes. *Comm. Math. Phys.*, 200(2):421–444, 1999.
- [16] Andrea Malchiodi and David Ruiz. New improved Moser-Trudinger inequalities and singular Liouville equations on compact surfaces. *Geom. Funct. Anal.*, 21(5):1196–1217, 2011.
- [17] G. Mancini. Onofri-type inequalities for singular liouville equations. *Journal of Geometric Analysis*, 2015.
- [18] J. Moser. A sharp form of an inequality by N. Trudinger. *Indiana Univ. Math. J.*, 20:1077–1092, 1970/71.
- [19] J. Moser. On a nonlinear problem in differential geometry. In *Dynamical systems (Proc. Sympos., Univ. Bahia, Salvador, 1971)*, pages 273–280. Academic Press, New York, 1973.
- [20] E. Onofri. On the positivity of the effective action in a theory of random surfaces. *Comm. Math. Phys.*, 86(3):321–326, 1982.
- [21] B. Osgood, R. Phillips, and P. Sarnak. Compact isospectral sets of surfaces. *J. Funct. Anal.*, 80(1):212–234, 1988.
- [22] B. Osgood, R. Phillips, and P. Sarnak. Extremals of determinants of Laplacians. *J. Funct. Anal.*, 80(1):148–211, 1988.
- [23] A. M. Polyakov. Quantum geometry of bosonic strings. *Phys. Lett. B*, 103(3):207–210, 1981.
- [24] A. M. Polyakov. Quantum geometry of fermionic strings. *Phys. Lett. B*, 103(3):211–213, 1981.

- [25] Michael Struwe. A flow approach to Nirenberg's problem. *Duke Math. J.*, 128(1):19–64, 2005.
- [26] Gabriella Tarantello. *Selfdual gauge field vortices*. Progress in Nonlinear Differential Equations and their Applications, 72. Birkhäuser Boston, Inc., Boston, MA, 2008. An analytical approach.
- [27] Marc Troyanov. Prescribing curvature on compact surfaces with conical singularities. *Trans. Amer. Math. Soc.*, 324(2):793–821, 1991.