

Self-Adjoint Extensions with Friedrichs Lower Bound

Matteo Gallone · Alessandro Michelangeli

Abstract

We produce a simple criterion and a constructive recipe to identify those self-adjoint extensions of a lower semi-bounded symmetric operator on Hilbert space which have the same lower bound as the Friedrichs extension. Applications of this abstract result to a few instructive examples are then discussed.

Keywords Lower semi-bounded symmetric operators · Self-adjoint extensions · Friedrichs extensions

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1 Motivation

We start with a familiar example. In the Hilbert space $\mathcal{H} = L^2(0, 1)$ let us consider the densely defined, closed, and symmetric operator

$$S = -\frac{d^2}{dx^2}, \quad \mathcal{D}(S) = H_0^2(0, 1) = \left\{ f \in H^2(0, 1) \mid \begin{array}{l} f(0) = 0 = f(1) \\ f'(0) = 0 = f'(1) \end{array} \right\}. \quad (1.1)$$

S is actually the operator closure of the negative Laplacian defined on $C_0^\infty(0, 1)$. Here and in the following $\mathcal{D}(R)$ denotes the domain of the operator R acting on \mathcal{H} , and if R is symmetric we denote by

$$m(R) := \inf_{\substack{f \in \mathcal{D}(R) \\ f \neq 0}} \frac{\langle f, Rf \rangle}{\|f\|_2^2} \in [-\infty, +\infty) \quad (1.2)$$

the largest lower bound of R . When $m(R) > -\infty$ one says that R is semi-bounded from below. Poincaré inequality implies that S is semi-bounded from below with

$$m(S) = \pi^2. \quad (1.3)$$

Now, S is symmetric but not self-adjoint, for

$$S^* = -\frac{d^2}{dx^2}, \quad \mathcal{D}(S^*) = H^2(0, 1). \quad (1.4)$$

Thus, S admits a multiplicity (in fact, a four-real-parameter family) of distinct self-adjoint extensions, which are all restrictions of S^* . Among them, the Friedrichs extension S_F is the one with domain

$$\mathcal{D}(S_F) = H^2(0, 1) \cap H_0^1(0, 1) = \{f \in H^2(0, 1) \mid f(0) = 0 = f(1)\}, \quad (1.5)$$

namely the Dirichlet (negative) Laplacian. Let us recall that abstractly speaking the Friedrichs extension of a lower semi-bounded symmetric operator S is the *only* self-adjoint extension of S with the property

$$\mathcal{D}(S_F) \subset \mathcal{D}[S], \quad (1.6)$$

that is, with operator domain contained in the form domain of S . Here and in the following $\mathcal{D}[R]$ denotes the form domain of a lower semi-bounded symmetric operator R , or also of a self-adjoint operator R (see, e.g., [11, Chapt. 10]); in the present case

$$\mathcal{D}[S] = \overline{\mathcal{D}(S)}^{\|\cdot\|_{H^1}} = H_0^1(0, 1), \quad (1.7)$$

and obviously $\mathcal{D}[S] = \mathcal{D}[S_F]$. It is also a general property of the Friedrichs extension the fact that $S_F \geq \tilde{S}$ for any other $\tilde{S} = \tilde{S}^* \supset S$, namely S_F is the largest of all the self-adjoint extensions of S in the sense of operator form ordering.

As well known, as follows from (1.5), S_F is diagonalizable over an orthonormal basis $\{\sqrt{2} \sin n\pi x \mid n \in \mathbb{N}\}$ of eigenfunctions, with (simple and pure point) spectrum

$$\sigma(S_F) = \{n^2\pi^2 \mid n \in \mathbb{N}\} \quad (1.8)$$

Thus,

$$m(S_F) = \pi^2 = m(S), \quad (1.9)$$

which actually expresses a completely general fact: the Friedrichs extension of a lower semi-bounded operator preserves the lower bound. Whereas self-adjoint extensions of S cannot increase the lower bound, in general they decrease it. For instance, in the example under consideration, the extension S_P with periodic boundary conditions, namely with domain

$$\mathcal{D}(S_P) = \{f \in H^2(0, 1) \mid f(1) = f(0), f'(1) = f'(0)\}, \quad (1.10)$$

has spectrum

$$\sigma(S_P) = \{n^2\pi^2 \mid n \in \mathbb{N}_0\}, \quad \text{whence} \quad m(S_P) = 0. \quad (1.11)$$

Yet, the extension S_A with anti-periodic boundary conditions, namely with

$$\mathcal{D}(S_A) = \{f \in H^2(0, 1) \mid f(1) = -f(0), f'(1) = -f'(0)\}, \quad (1.12)$$

does preserve the lower bound of S . Indeed,

$$\sigma(S_A) = \{(2n+1)^2\pi^2 \mid n \in \mathbb{N}_0\}, \quad \text{whence} \quad m(S_P) = \pi^2. \quad (1.13)$$

This occurrence is well known: a lower semi-bounded symmetric operator may admit self-adjoint extensions other than the Friedrichs, with the same bottom of the Friedrichs spectrum. Actually this is not typical of symmetric operators with deficiency index 2 only, as was the case for S here. In Sect. 3 also examples with deficiency indices 1 will be recalled and discussed. By standard direct sum, these examples also cover the case of infinite deficiency indices.

Now, while the possibility of non-Friedrichs self-adjoint extensions with the same Friedrichs lower bound is folk knowledge, we are not aware of an explicit operator-theoretic explanation of this phenomenon, nor of a characterisation in terms of transparent conditions which, once they are met, allow to construct all extensions with such a feature.

In this note we present a simple criterion and a constructive recipe to identify those self-adjoint extensions of a lower semi-bounded symmetric operator on Hilbert space which have the same lower bound as the Friedrichs extension. The abstract main results, Theorems 2.4, 2.5, and 2.6 below, are discussed in Sect. 2, and illustrative concrete examples where such results can be applied to are then presented in Sect. 3.

2 Abstract Results

Let \mathcal{H} be a Hilbert space (over \mathbb{R} or \mathbb{C} , with scalar product $\langle \cdot, \cdot \rangle$ anti-linear in the first entry, and with norm $\|\cdot\|$) and let S be a densely defined, symmetric, semi-bounded

operator on \mathcal{H} with lower bound $m(S)$. S is not necessarily closed. For clarity of the presentation we shall assume non-restrictively $m(S) > 0$. This implies that S_F^{-1} is everywhere defined and bounded on \mathcal{H} . It will be clear both from this abstract discussion and from the applications in Sect. 3 that the case of general finite $m(S)$ can be covered by suitably shifting S to $S - \mu \mathbb{1}$ with $\mu < m(S)$.

Unless such S is already essentially self-adjoint, it admits non-trivial self-adjoint extensions. In this case $\ker S^*$, the deficiency space for S , is non-trivial either. Standard extension schemes produce convenient classifications of the whole family of extensions. It can be shown within the modern theory of boundary triplets [2], or equivalently the classical ‘universal’ parametrization by Grubb [7], and in fact the very original extension theory by Kreĭn [9], Višik [12], and Birman [3], that the extensions of S can be labelled as follows.

Theorem 2.1 *There is a one-to-one correspondence between the family of all self-adjoint extensions of S on \mathcal{H} and the family of the self-adjoint operators on Hilbert subspaces of $\ker S^*$.*

(i) *If T is any such operator, in the correspondence $T \leftrightarrow S_T$ each self-adjoint extension S_T of S is given by*

$$\begin{aligned} S_T &= S^* \upharpoonright \mathcal{D}(S_T) \\ \mathcal{D}(S_T) &= \left\{ f + S_F^{-1}(Tv + w) + v \left| \begin{array}{l} f \in \mathcal{D}(\bar{S}), \quad v \in \mathcal{D}(T) \\ w \in \ker S^* \cap \mathcal{D}(T)^\perp \end{array} \right. \right\}. \end{aligned} \quad (2.1)$$

(ii) *If S_T is a semi-bounded (not necessarily positive) self-adjoint extension of S , then*

$$\mathcal{D}[T] = \mathcal{D}[S_T] \cap \ker S^* \quad (2.2)$$

and

$$\begin{aligned} \mathcal{D}[S_T] &= \mathcal{D}[S_F] \dot{+} \mathcal{D}[T] \\ S_T[f + v, f' + v'] &= S_F[f, f'] + T[v, v'] \\ \forall f, f' \in \mathcal{D}[S_F], \forall v, v' \in \mathcal{D}[T]. \end{aligned} \quad (2.3)$$

As a consequence,

$$S_{T_1} \geq S_{T_2} \quad \Leftrightarrow \quad T_1 \geq T_2. \quad (2.4)$$

(iii) *If $m(T) > -m(S)$, then*

$$m(T) \geq m(S_T) \geq \frac{m(S) m(T)}{m(S) + m(T)}. \quad (2.5)$$

Theorem 2.1 collects results that are proved, e.g., in [8, Chapt. 13], [11, Chapt. 14], and [5, Sect. 3].

For convenience, let us denote by $\mathcal{S}(\mathcal{K})$ the collection of all self-adjoint operators defined in Hilbert subspaces of a given Hilbert space \mathcal{K} : Theorem 2.1 states that the self-adjoint extensions of S are all of the form S_T for some $T \in \mathcal{S}(\ker S^*)$.

The Friedrichs extension of S can be expressed in terms of the classical decomposition formula (see, e.g., [5, Sect. 2.2])

$$\mathcal{D}(S_F) = \mathcal{D}(\bar{S}) \dot{+} S_F^{-1} \ker S^*. \tag{2.6}$$

Therefore, S_F is recovered from the general parametrisation (2.1) or (2.3) with the choice $\mathcal{D}[T] = \{0\}$ (thus, formally, “ $T = \infty$ ”).

An ancillary result that tends to be somehow less highlighted, but which is most relevant for our discussion, is the following.

Theorem 2.2 *If, with respect to the notation of (2.1), S_T is a self-adjoint extension of S , and if $\mu < m(S)$, then*

$$\begin{aligned} \langle g, S_T g \rangle &\geq \mu \|g\|^2 \quad \forall g \in \mathcal{D}(S_T) \\ &\Updownarrow \\ \langle v, T v \rangle &\geq \mu \|v\|^2 + \mu^2 \langle v, (S_F - \mu \mathbb{1})^{-1} v \rangle \quad \forall v \in \mathcal{D}(T). \end{aligned} \tag{2.7}$$

As an immediate consequence, Theorem 2.2 reproduces the inequality $m(T) \geq m(S_T)$ for any semi-bounded S_T and shows, in particular, that positivity or strict positivity of the bottom of S_T is equivalent to the same property for T , that is,

$$\begin{aligned} m(S_T) \geq 0 &\Leftrightarrow m(T) \geq 0 \\ m(S_T) > 0 &\Leftrightarrow m(T) > 0. \end{aligned} \tag{2.8}$$

To make this presentation self-contained, and for later convenience, let us deduce Theorem 2.2 from Theorem 2.1. To this aim, let us first single out a simple operator-theoretic property.

Lemma 2.3 *If A is a self-adjoint operator on a Hilbert space \mathcal{H} with positive bottom ($m(A) > 0$), then*

$$\sup_{f \in \mathcal{D}(A)} \frac{|\langle f, h \rangle|^2}{\langle f, Af \rangle} = \langle h, A^{-1} h \rangle \quad \forall h \in \mathcal{H}.$$

Proof Setting $g := A^{1/2} f$ one has

$$\sup_{f \in \mathcal{D}(A)} \frac{|\langle f, h \rangle|^2}{\langle f, Af \rangle} = \sup_{g \in \mathcal{H}} \frac{|A^{-1/2} g, h|^2}{\|g\|^2} = \sup_{\|g\|=1} |\langle g, A^{-1/2} h \rangle|^2$$

and since $|\langle g, A^{-1/2} h \rangle|$ attains its maximum for $g = A^{-1/2} h / \|A^{-1/2} h\|$, the conclusion then follows. □

Proof of Theorem 2.2 For generic $f \in \mathcal{D}(S_F)$ and $v \in \mathcal{D}(T)$, one has $g := f + v \in \mathcal{D}(S_T)$ and

$$S_T[g] = \langle f, S_F f \rangle + \langle v, T v \rangle.$$

Thus, $S_T \geq \mu \mathbb{1}$ is tantamount as requiring for all such g 's that

$$\langle f, S_F f \rangle + \langle v, T v \rangle \geq \mu (\langle f, f \rangle + \langle f, v \rangle + \langle v, f \rangle + \langle v, v \rangle)$$

whence also, replacing $f \mapsto \lambda f$, $v \mapsto \gamma v$,

$$\begin{aligned} & (\langle f, S_F f \rangle - \mu \|f\|^2) |\lambda|^2 - \mu \langle f, v \rangle \bar{\lambda} \gamma - \mu \langle v, f \rangle \lambda \bar{\gamma} \\ & + (\langle v, T v \rangle - \mu \|v\|^2) |\gamma|^2 \geq 0 \quad \forall \lambda, \gamma \in \mathbb{C}. \end{aligned}$$

Since $\mu < m(S)$, and hence $\langle f, S_F f \rangle - \mu \|f\|^2 > 0$, last inequality holds true if and only if

$$\mu^2 |\langle f, v \rangle|^2 \leq (\langle v, T v \rangle - \mu \|v\|^2) (\langle f, S_F f \rangle - \mu \|f\|^2) \quad (*)$$

for arbitrary $f \in \mathcal{D}(S_F)$ and $v \in \mathcal{D}(T)$. By re-writing (*) as

$$\langle v, T v \rangle - \mu \|v\|^2 \geq \mu^2 \frac{|\langle f, v \rangle|^2}{\langle f, (S_F - \mu \mathbb{1}) f \rangle}$$

and by the fact that the above inequality is valid for arbitrary $f \in \mathcal{D}(S_F)$ and hence holds true also when the supremum over such f 's is taken, Lemma 2.3 then yields

$$\langle v, T v \rangle - \mu \|v\|^2 \geq \mu^2 \langle v, (S_F - \mu \mathbb{1})^{-1} v \rangle,$$

which completes the proof. \square

With these abstract results at hand, let us now turn to the identification of the non-Friedrichs extensions with the same Friedrichs lower bound.

It is worth observing that inequality (2.5) is not informative in this respect: indeed, owing to (2.5), a *sufficient* condition for the bottom of S_T to equal the bottom of S_F would be to impose $m(S)m(T)/(m(S) + m(T)) \geq m(S)$, but such inequality is only satisfied, in the form of an identity, when $m(T) = \infty$, therefore the above sufficient condition only selects $S_T = S_F$, the Friedrichs extension.

We rather focus on (2.7) from Theorem 2.2. There, the operator $S_F - \mu \mathbb{1}$ is invertible with everywhere bounded inverse on the whole \mathcal{H} : indeed, $\mu < m(S)$ and then $m(S_F - \mu \mathbb{1}) > 0$. Instead, $S_F - m(S)\mathbb{1}$ fails to be invertible on \mathcal{H} , because its bottom is by construction equal to zero.

The informal idea now is that even if $(S_F - m(S)\mathbb{1})^{-1}$ cannot be defined as a bounded operator *on the whole* \mathcal{H} , yet it makes sense on $\text{ran}(S_F - m(S)\mathbb{1})$, and if it happens that the latter space has a non-trivial intersection with $\ker S^*$, then there are

non-zero vectors $v \in \text{ran}(S_F - \mathfrak{m}(S)\mathbb{1}) \cap \ker S^*$ on which $\langle v, (S_F - \mathfrak{m}(S)\mathbb{1})^{-1}v \rangle$ is unambiguously defined and hence the right-hand side of the second expression in (2.7) is meaningful also when $\mu = \mathfrak{m}(S)$. Moreover, if on such v 's one can define an operator $T \in \mathcal{S}(\ker S^*)$ satisfying (2.7) when $\mu = \mathfrak{m}(S)$, then by suitably exploiting the limit $\mu \uparrow \mathfrak{m}(S)$ this should give a characterisation of the condition $\mathfrak{m}(S_T) \geq \mathfrak{m}(S)$, which is equivalent to $\mathfrak{m}(S_T) = \mathfrak{m}(S)$, as $S_F \geq S_T$, the Friedrichs extension is the largest of all self-adjoint extensions of S .

By elaborating on such idea we finally come to our main results, Theorems 2.4, 2.5, and 2.6 below.

Clearly, underlying (2.7) is the quadratic form language, so the actual operator to possibly invert in some subspace of $\ker S^*$ is rather $(S_F - \mathfrak{m}(S)\mathbb{1})^{1/2}$, a positive self-adjoint operator with zero lower bound.

In this respect, as S_F is self-adjoint on \mathcal{H} , and so is $S_F - \mathfrak{m}(S)\mathbb{1}$ with lower bound zero, then upon decomposing

$$\mathcal{H} = \overline{\text{ran}(S_F - \mathfrak{m}(S)\mathbb{1})} \oplus \ker(S_F - \mathfrak{m}(S)\mathbb{1})$$

the negative powers $(S_F - \mathfrak{m}(S)\mathbb{1})^{-\delta}$, $\delta > 0$, are naturally defined as self-adjoint operators on the Hilbert subspace $\overline{\text{ran}(S_F - \mathfrak{m}(S)\mathbb{1})}$, or also on the whole \mathcal{H} upon extension by zero on $\ker(S_F - \mathfrak{m}(S)\mathbb{1})$.

In the first statement we *characterise* the occurrence of non-Friedrichs extensions with the same Friedrichs lower bound.

Theorem 2.4 *Let S be a densely defined and symmetric operator on a given Hilbert space \mathcal{H} with lower bound $\mathfrak{m}(S) > 0$. Necessary and sufficient condition for S to admit self-adjoint extensions other than the Friedrichs extensions and with the same lower bound $\mathfrak{m}(S)$ is that*

$$\text{ran}(S_F - \mathfrak{m}(S)\mathbb{1})^{1/2} \cap \ker S^* \neq \{0\}. \quad (2.9)$$

In the applications both $\ker S^*$ and $\text{ran}(S_F - \mathfrak{m}(S)\mathbb{1})^{1/2}$ are in general spaces that one can qualify rather explicitly. Thus, condition (2.9) is practically manageable and qualifies the operator-theoretic *mechanism* for non-Friedrichs extensions with the Friedrichs lower bound. In Sect. 3 we shall give examples of that.

Our next result concerns the actual *recipe* to construct such extensions, when (2.9) is matched, thus in practice how to identify the corresponding extension parameters T in $\mathcal{S}(\ker S^*)$. We shall use the customary notation of square brackets for the domain $\mathcal{D}[q]$ of a quadratic form q on \mathcal{H} and for the evaluation $q[v]$ on elements of its domain; as usual, we shall denote by $q[v_1, v_2]$ the evaluation of the corresponding sesquilinear form defined by polarisation.

Theorem 2.5 *Same assumptions as in Theorem 2.4, and assume further that condition (2.9) is satisfied.*

(i) *The expression*

$$\begin{aligned} \mathcal{D}[q] &:= \text{ran}(S_F - \mathfrak{m}(S)\mathbb{1})^{\frac{1}{2}} \cap \ker S^* \\ q[v] &:= \mathfrak{m}(S)\|v\|^2 + \mathfrak{m}(S)^2 \|(S_F - \mathfrak{m}(S)\mathbb{1})^{-\frac{1}{2}}v\|^2. \end{aligned} \quad (2.10)$$

defines a symmetric, closed, and strictly positive quadratic form q .

- (ii) Let T_q be the operator on the Hilbert subspace $\overline{\mathcal{D}[q]}$ uniquely associated with q . Then $T_q \in \mathcal{S}(\ker S^*)$.
- (iii) For any $T \in \mathcal{S}(\ker S^*)$ with $T \geq T_q$, the corresponding self-adjoint extension S_T of S (Theorem 2.1) has the property

$$\mathfrak{m}(S_T) = \mathfrak{m}(S). \quad (2.11)$$

- (iv) Any self-adjoint extension S_T of S satisfying (2.11) corresponds to an extension parameter $T \in \mathcal{S}(\ker S^*)$ with $T \geq T_q$.

By definition, in (2.10) the vector $u_\circ = (S_F - \mathfrak{m}(S)\mathbb{1})^{-\frac{1}{2}}v$ is the minimal norm solution $u = u_\circ$ to $(S_F - \mathfrak{m}(S)\mathbb{1})^{\frac{1}{2}}u = v$.

In view of the general classification of Theorem 2.1, the above results admit a natural corollary that it is worth stating as a separate theorem. It is convenient to introduce the meaningful terminology ‘*top extensions*’ for all those S_T ’s with $\mathfrak{m}(S_T) = \mathfrak{m}(S)$ (in particular, S_F is a top extension), and ‘*least-top extension*’ for the extension $S_{LT} := S_{T_q}$.

Theorem 2.6 *Same assumptions as in Theorem 2.4. Each top extension S^{top} of S satisfies*

$$S_F \geq S^{\text{top}} \geq S_{LT} \quad (2.12)$$

in the sense of operator form ordering. Each such extension is of the form $S^{\text{top}} = S_T$ for some $T \in \mathcal{S}(\ker S^*)$ with $T \geq T_q$, where T_q is qualified in Theorem 2.5(ii), and they are all ordered with T in the sense of (2.4).

Proof (Proof of Theorem 2.4, necessity part)

Let S_T be a self-adjoint extension of S , labelled by some $T \in \mathcal{S}(\ker S^*)$, with the property $\mathfrak{m}(S_T) = \mathfrak{m}(S_F)$ and $S_T \neq S_F$. Let $(\mu_n)_{n \in \mathbb{N}}$ be an increasing sequence of real numbers such that $\mu_n < \mathfrak{m}(S) \forall n$ and $\mu_n \rightarrow \mathfrak{m}(S)$ as $n \rightarrow \infty$. Since $S_T \geq \mu_n \mathbb{1}$, Theorem 2.2 implies

$$\langle v, T v \rangle \geq \mu_n \|v\|^2 + \mu_n^2 \|(S_F - \mu_n \mathbb{1})^{-\frac{1}{2}}v\|^2$$

for every $v \in \mathcal{D}(T)$, whence

$$\limsup_{n \rightarrow \infty} \|(S_F - \mu_n \mathbb{1})^{-\frac{1}{2}}v\|^2 < +\infty.$$

In fact, for each v the sequence of square norms $\|(S_F - \mu_n \mathbb{1})^{-\frac{1}{2}} v\|^2$ is monotone increasing. For, if $m(S) > \mu' > \mu$, then

$$\begin{aligned} & \|(S_F - \mu' \mathbb{1})^{-\frac{1}{2}} v\|^2 - \|(S_F - \mu \mathbb{1})^{-\frac{1}{2}} v\|^2 \\ &= \int_{[m(S), +\infty)} \left(\frac{1}{\lambda - \mu'} - \frac{1}{\lambda - \mu} \right) dv_v(\lambda) \geq 0, \end{aligned}$$

where ν_v is the scalar spectral measure of the self-adjoint operator S_F relative to the vector v . Therefore,

$$\exists \lim_{n \rightarrow \infty} \|(S_F - \mu_n \mathbb{1})^{-\frac{1}{2}} v\|^2 < +\infty.$$

As the latter conclusion is tantamount as

$$\exists \lim_{n \rightarrow \infty} \int_{[m(S), +\infty)} \frac{1}{\lambda - \mu_n} dv_v(\lambda) < +\infty,$$

then by monotone convergence the function $\lambda \mapsto (\lambda - m(S))^{-1}$ is ν_v -summable. Thus, $\|(S_F - m(S)\mathbb{1})^{-\frac{1}{2}} v\|^2 < +\infty$, whence $v \in \text{ran}(S_F - m(S)\mathbb{1})^{\frac{1}{2}}$.

On the other hand, since by assumption S_T is a self-adjoint extension of S distinct from S_F , then by definition of extension parameter T one has that $\mathcal{D}(T)$ is a non-trivial subspace of $\ker S^*$. Summarising,

$$\mathcal{D}(T) \subset \text{ran}(S_F - m(S)\mathbb{1})^{\frac{1}{2}} \cap \ker S^*$$

and therefore $\text{ran}(S_F - m(S)\mathbb{1})^{\frac{1}{2}} \cap \ker S^*$ is non-trivial. □

Corollary 2.7 *As a consequence of the necessity statement of Theorem 2.4, for each self-adjoint extension S_T of S with $m(S_T) = m(S)$ one has*

$$\begin{aligned} & \mathcal{D}(T) \subset \text{ran}(S_F - m(S)\mathbb{1})^{\frac{1}{2}} \cap \ker S^* \quad \text{and} \\ & \langle v, Tv \rangle \geq m(S)\|v\|^2 + m(S)^2 \|(S_F - m(S)\mathbb{1})^{-\frac{1}{2}} v\|^2 \quad \forall v \in \mathcal{D}(T). \end{aligned} \tag{2.13}$$

Proof The inclusion for $\mathcal{D}(T)$ was already proved. Next, as a follow-up of the reasoning of the previous proof, let us observe that for each $v \in \mathcal{D}(T)$ one has

$$\lim_{n \rightarrow \infty} \|(S_F - \mu_n \mathbb{1})^{-\frac{1}{2}} v\|^2 = \|(S_F - m(S)\mathbb{1})^{-\frac{1}{2}} v\|^2.$$

Indeed,

$$\begin{aligned} & \|(S_F - m(S)\mathbb{1})^{-\frac{1}{2}} v\|^2 - \|(S_F - \mu_n \mathbb{1})^{-\frac{1}{2}} v\|^2 \\ &= \int_{[m(S), +\infty)} \left(\frac{1}{\lambda - m(S)} - \frac{1}{\lambda - \mu_n} \right) dv_v(\lambda) \xrightarrow{n \rightarrow \infty} 0 \end{aligned}$$

by dominated convergence. Therefore, one can take the limit $n \rightarrow \infty$ in the inequality

$$\langle v, Tv \rangle \geq \mu_n \|v\|^2 + \mu_n^2 \|(S_F - \mu_n \mathbb{1})^{-\frac{1}{2}} v\|^2$$

thus obtaining the second line of (2.13). \square

Proof (Proof of Theorem 2.5 and of Theorem 2.4, sufficiency part)

(i) The fact that (2.10) defines a symmetric quadratic form with strictly positive lower bound is obvious. As for q being closed, let us show that if $(v_n)_{n \in \mathbb{N}}$ is a sequence in $\mathcal{D}[q]$ with $v_n \rightarrow v$ and $q[v_n - v_m] \rightarrow 0$ as $n, m \rightarrow \infty$, then $v \in \mathcal{D}[q]$ and $q[v_n - v] \rightarrow 0$. This is indeed equivalent to saying that q is closed (see, e.g., [11, Prop. 10.1]). Now, the above assumption on $(v_n)_{n \in \mathbb{N}}$ implies

$$\begin{aligned} v_n &\rightarrow v \\ (S_F - m(S)\mathbb{1})^{-\frac{1}{2}} v_n &\rightarrow u \end{aligned}$$

for some $v \in \mathcal{H}$. As $(S_F - m(S)\mathbb{1})^{-\frac{1}{2}}$ is self-adjoint and hence closed, this implies

$$\begin{aligned} v &\in \mathcal{D}((S_F - m(S)\mathbb{1})^{-\frac{1}{2}}) = \text{ran}((S_F - m(S)\mathbb{1})^{\frac{1}{2}}) \\ u &= (S_F - m(S)\mathbb{1})^{-\frac{1}{2}} v. \end{aligned}$$

A first conclusion, since $\ker S^*$ is closed in \mathcal{H} and hence $v \in \ker S^*$ as well, is that $v \in \text{ran}((S_F - m(S)\mathbb{1})^{\frac{1}{2}}) \cap \ker S^* = \mathcal{D}[q]$. As further conclusion, since $v_n \rightarrow v$ and $(S_F - m(S)\mathbb{1})^{-\frac{1}{2}} v_n \rightarrow (S_F - m(S)\mathbb{1})^{-\frac{1}{2}} v$ in \mathcal{H} , one has $q[v_n - v] \rightarrow 0$. Part (i) of Theorem 2.5 is thus proved.

(ii) As q is densely defined in the Hilbert subspace $\overline{\mathcal{D}[q]}$, and it is symmetric, closed, and semi-bounded from below, then q uniquely identifies a self-adjoint operator T_q on $\overline{\mathcal{D}[q]}$ defined by

$$\begin{aligned} \mathcal{D}(T_q) &:= \{v \in \mathcal{D}[q] \mid \exists z_v \in \mathcal{H} \text{ with } \langle u, z_v \rangle = q[u, v] \ \forall u \in \mathcal{D}[q]\} \\ T_q v &:= z_v \end{aligned}$$

(see, e.g., [11, Theorem 10.7]). Since $\mathcal{D}[q] \subset \ker S^*$ and $\ker S^*$ is closed in \mathcal{H} , then $\overline{\mathcal{D}[q]} \subset \ker S^*$, thus proving that $T_q \in \mathcal{S}(\ker S^*)$. This establishes part (ii) of Theorem 2.5.

(iii) Let $T \in \mathcal{S}(\ker S^*)$ with $T \geq T_q$. This means that $\mathcal{D}(T) \subset \mathcal{D}(T_q)$ and

$$\langle v, Tv \rangle \geq \langle v, T_q v \rangle = q[v] = m(S)\|v\|^2 + m(S)^2 \|(S_F - m(S)\mathbb{1})^{-\frac{1}{2}} v\|^2$$

for every $v \in \mathcal{D}(T)$. Consider now an arbitrary $\mu < m(S)$. With the very same argument used in the proof of the necessity part of Theorem 2.5 one sees that

$$\|(S_F - m(S)\mathbb{1})^{-\frac{1}{2}} v\|^2 > \|(S_F - \mu\mathbb{1})^{-\frac{1}{2}} v\|^2,$$

whence

$$\langle v, Tv \rangle > \mu \|v\|^2 + \mu^2 \|(S_F - \mu \mathbb{1})^{-\frac{1}{2}} v\|^2$$

for all $v \in \mathcal{D}(T)$. Owing to Theorem 2.2, the self-adjoint extension S_T of S parametrised by the considered T is such that $S_T \geq \mu \mathbb{1}$. By the arbitrariness of μ , one concludes that $S_T \geq m(S)\mathbb{1}$, whence $m(S_T) = m(S)$. Unless “ $T = \infty$ ” (in the sense $\mathcal{D}[T] = \{0\}$), all other choices for T identifies non-Friedrichs extensions. This completes the proof of part (iii) of Theorem 2.5. At the same time, this proves that assumption (2.9) in Theorem 2.4 allows one to construct non-Friedrichs extensions with the same Friedrichs lower bound. Thus also the sufficiency statement of Theorem 2.4 is established.

(iv) Last, let S_T be a self-adjoint extension of S with $m(S_T) = m(S)$. The necessity statement of Theorem 2.4 implies that the intersection $\text{ran}(S_F - m(S)\mathbb{1})^{\frac{1}{2}} \cap \ker S^*$ is non-trivial, so one can define the form q and the operator $T_q \in \mathcal{S}(\ker S^*)$ as in parts (i) and (ii) of Theorem 2.5. Owing to Corollary 2.7,

$$\begin{aligned} \mathcal{D}[T] &\subset \mathcal{D}(q) \\ T[v] &= q[v] \quad \forall v \in \mathcal{D}[T]. \end{aligned}$$

This means precisely that $T \geq T_q$. □

3 Applications

Let us discuss now a few instructive examples of application of Theorems 2.4 and 2.5. We present examples with deficiency index equal to two (Sect. 3.1) and one (Sects. 3.2, 3.3, 3.4), and we cover also the edge case in which the sub-family of top extensions consists of the Friedrichs extension only (Sect. 3.4).

3.1 Schrödinger Quantum Particle on an Interval

Let us revisit in more systematic terms the example presented in Sect. 1. The operator S has deficiency index equal to 2, and explicitly

$$\ker S^* = \text{span}\{\mathbf{1}, x\}. \quad (3.1)$$

The operator $S_F - \pi^2 \mathbb{1}$ fails to be invertible on the whole $\mathcal{H} = L^2(0, 1)$ because it has a non-trivial kernel:

$$\begin{aligned} \ker(S_F - \pi^2 \mathbb{1}) &= \text{span}\{\sin \pi x\} \\ \text{ran}(S_F - \pi^2 \mathbb{1}) &= \text{span}\{\sin n\pi x \mid n \in \mathbb{N}, n \geq 2\}. \end{aligned} \quad (3.2)$$

As $(S_F - \pi^2 \mathbb{1})$ is diagonalised as above over an orthonormal basis of eigenfunctions, its powers $(S_F - \pi^2 \mathbb{1})^\delta$ and $(S_F - \pi^2 \mathbb{1})^{-\delta}$, with $\delta > 0$, are qualified by their action on the same basis of eigenfunctions, with eigenvalues given by the corresponding powers of the eigenvalues of $(S_F - \pi^2 \mathbb{1})$; the negative powers are clearly only defined on the Hilbert subspace $\text{ran}(S_F - \pi^2 \mathbb{1}) = \{\sin \pi x\}^\perp$. Therefore,

$$\text{ran}(S_F - \pi^2 \mathbb{1})^{\frac{1}{2}} = \text{ran}(S_F - \pi^2 \mathbb{1}) = \text{span}\{\sin n\pi x \mid n \in \mathbb{N}, n \geq 2\}. \quad (3.3)$$

Lemma 3.1 *One has*

$$V := \text{ran}(S_F - \pi^2 \mathbb{1})^{\frac{1}{2}} \cap \ker S^* = \text{span}\{\mathbf{1} - 2x\} \quad (3.4)$$

and

$$(S_F - \pi^2 \mathbb{1})^{-1}(\mathbf{1} - 2x) = \pi^{-2}(\cos \pi x - \mathbf{1} + 2x). \quad (3.5)$$

Proof In order for a generic element $a\mathbf{1} + bx \in \ker S^*$, with $a, b \in \mathbb{C}$, to belong to $\text{ran}(S_F - \pi^2 \mathbb{1})^{\frac{1}{2}}$, owing to (3.2)–(3.3) it must be

$$0 = \int_0^1 (a + bx) \sin \pi x \, dx = \pi^{-1}(2a + b),$$

whence $b = -2a$. Thus $g \in V$ implies $g = a(\mathbf{1} - 2x)$ for some $a \in \mathbb{C}$. Next, one has to check that $\mathbf{1} - 2x \in \text{ran}(S_F - \pi^2 \mathbb{1})^{\frac{1}{2}}$. This is the same as $\mathbf{1} - 2x \in \text{ran}(S_F - \pi^2 \mathbb{1})$, that is,

$$\mathbf{1} - 2x = (S_F - \pi^2 \mathbb{1})u \quad \text{for some } u \in \mathcal{D}(S_F).$$

This is equivalent to saying that u is the *minimal norm solution* to the boundary value problem

$$\begin{cases} -u'' - \pi^2 u = 1 - 2x \\ u(0) = 0 = u(1). \end{cases}$$

By standard ODE methods one finds that the general solution is

$$u_{\text{gen}}(x) = \pi^{-2}(\cos \pi x - 1 + 2x) + B \sin \pi x, \quad B \in \mathbb{C},$$

thus the minimal norm solution is the one with $B = 0$. This proves that the function $u_\circ := \pi^{-2}(\cos \pi x - \mathbf{1} + 2x) \in \mathcal{D}(S_F)$ satisfies $(S_F - \pi^2 \mathbb{1})u_\circ = \mathbf{1} - 2x$, thus completing the proof of (3.4) and (3.5). \square

As the intersection space (3.4) is non-trivial, Theorem 2.4 ensures that S admits non-Friedrichs self-adjoint extensions with the same Friedrichs lower bound. This is

consistent with what discussed in the introduction: $m(S_F) = m(S_A)$, Friedrichs and anti-periodic extension have the same lower bound.

It is instructive to apply the constructive recipe of Theorem 2.5 so as to identify all such extensions. With the notation therein,

$$\mathcal{D}[q] = \mathcal{D}(T_q) = V = \text{span}\{\mathbf{1} - 2x\}, \tag{3.6}$$

thus T_q is an operator of multiplication by some real number t_q ,

$$T_q(\mathbf{1} - 2x) = t_q(\mathbf{1} - 2x). \tag{3.7}$$

Since

$$\begin{aligned} \langle (\mathbf{1} - 2x), T_q(\mathbf{1} - 2x) \rangle &= \pi^2 \|\mathbf{1} - 2x\|_2^2 + \pi^4 \|(S_F - \pi^2 \mathbb{1})^{-\frac{1}{2}}(\mathbf{1} - 2x)\|_2^2 \\ &= \pi^2 \|\mathbf{1} - 2x\|_2^2 + \pi^4 \langle (\mathbf{1} - 2x), (S_F - \pi^2 \mathbb{1})^{-1}(\mathbf{1} - 2x) \rangle \\ &= \pi^2 \|\mathbf{1} - 2x\|_2^2 + \pi^2 \langle (\mathbf{1} - 2x), \cos x - (\mathbf{1} - 2x) \rangle \\ &= \pi^2 \langle (\mathbf{1} - 2x), \cos x \rangle \\ &= 4 = 12 \|\mathbf{1} - 2x\|_2^2 \end{aligned}$$

(having used (3.5) in the third step and $\|\mathbf{1} - 2x\|_2^2 = \frac{1}{3}$ in the last step), then necessarily $t_q = 12$.

Theorem 2.5, in parts (iii) and (iv), then states that the self-adjoint extensions S_T of S with $m(S_T) = m(S_F)$ are those labelled by self-adjoint operators T with $T \geq T_q$. Such T 's, apart from the one parametrising the Friedrichs extension, are therefore such that

$$\begin{aligned} \mathcal{D}(T) &= V = \text{span}\{\mathbf{1} - 2x\} \\ T &\text{ is the multiplication by some } t \geq 12. \end{aligned} \tag{3.8}$$

Keeping into account, as is immediate to check, that

$$W := V^\perp \cap \ker S^* = \text{span}\{\mathbf{1}\}, \tag{3.9}$$

the extension S_T for each T satisfying (3.8) has domain given by formula (2.1) of Theorem 2.1, that is,

$$\begin{aligned} \mathcal{D}(S_T) &= \left\{ f + S_F^{-1}(Tv + w) + v \left| \begin{array}{l} f \in \mathcal{D}(\bar{S}), \\ v \in V, w \in W \end{array} \right. \right\} \\ &= \left\{ f + S_F^{-1}(t\alpha(\mathbf{1} - 2x) + \beta\mathbf{1}) + \alpha(\mathbf{1} - 2x) \left| \begin{array}{l} f \in H_0^2(0, 1) \\ \alpha, \beta \in \mathbb{C} \end{array} \right. \right\}. \end{aligned} \tag{3.10}$$

The action of the everywhere defined and bounded operator S_F^{-1} on the subspace $\ker S^* = \text{span}\{\mathbf{1}, x\}$ is easily computed by solving a boundary value problem completely analogous to the one considered in the proof of Lemma 3.1. The result (as

found, e.g., in [5, Eq. (91)]) is

$$S_F^{-1}(a\mathbf{1} + bx) = \left(\frac{a}{2} + \frac{b}{6}\right)x - \frac{a}{2}x^2 - \frac{b}{6}x^3. \quad (3.11)$$

Thus,

$$\mathcal{D}(S_T) = \left\{ f + \alpha\mathbf{1} + \left(\frac{t\alpha + 3\beta}{6} - 2\alpha\right)x - \frac{t\alpha + \beta}{2}x^2 + \frac{t\alpha}{3}x^3 \mid \begin{array}{l} f \in H_0^2(0, 1) \\ \alpha, \beta \in \mathbb{C} \end{array} \right\}. \quad (3.12)$$

Formula (3.12), for each fixed $t \geq 12$, identifies those self-adjoint extensions of S different from the Friedrichs extension, but with the same lower bound. In order to identify the boundary condition of self-adjointness satisfied by a generic element $g \in \mathcal{D}(S_T)$ for each extension of type (3.12), we compute the boundary values

$$\begin{aligned} g(0) &= \alpha & g'(0) &= \frac{t\alpha}{6} + \frac{\beta}{2} - 2\alpha \\ g(1) &= -\alpha & g'(1) &= \frac{t\alpha}{6} - \frac{\beta}{2} - 2\alpha \end{aligned} \quad (3.13)$$

and re-write

$$\begin{aligned} g(0) + g(1) &= 0 \\ g'(0) + g'(1) &= \frac{1}{3}(t - 12)g(0). \end{aligned} \quad (3.14)$$

It was indeed convenient to cast (3.13) in the form (3.14) because the latter can be more easily matched with the general conditions of self-adjointness of the extensions of S , as we shall now do.

We refer to the following very standard result, obtained for example by exploiting Theorem 2.1 for all possible extension parameters (see, e.g., [11, Example 14.10]), or equivalently by means of the alternative extension scheme a la von Neumann applied to S (see, e.g., [6, Sect. 6.2.3.1]).

Proposition 3.2 *The family of self-adjoint extensions on $L^2(0, 1)$ of the operator S defined in (1.1) consists of restrictions of S^* , and hence of operators of the form $-\frac{d^2}{dx^2}$, to domains of $H^2(0, 1)$ -functions g satisfying boundary conditions of one of the following four classes:*

$$g'(0) = b_1g(0) + cg(1), \quad g'(1) = -\bar{c}g(0) - b_2g(1), \quad (3.15)$$

$$g'(0) = b_1g(0) + \bar{c}g'(1), \quad g(1) = cg(0), \quad (3.16)$$

$$g'(1) = -b_1g(1), \quad g(0) = 0, \quad (3.17)$$

$$g(0) = 0 = g(1), \quad (3.18)$$

where $c \in \mathbb{C}$ and $b_1, b_2 \in \mathbb{R}$ and qualify each extension.

Remark 3.3 With reference to the general formula (2.1), extensions of type (3.15) correspond to the case in which $\dim \mathcal{D}(T) = 2$, extensions of type (3.16) or (3.17) correspond to $\dim \mathcal{D}(T) = 1$, and the extension of type (3.18) is the Friedrichs extension, $\dim \mathcal{D}(T) = 0$.

By direct comparison between (3.14) and (3.15)–(3.18) we see that (3.14) can only be of type (3.16) with

$$b_1 = \frac{1}{3}(t - 12) \quad \text{and} \quad c = -1. \quad (3.19)$$

We have thus proved the following.

Proposition 3.4 *The non-Friedrichs self-adjoint extensions on $L^2(0, 1)$ of the operator S defined in (1.1) which preserve the Friedrichs lower bound $\mathfrak{m}(S) = \pi^2$ are all those operators acting as $-\frac{d^2}{dx^2}$ on $H^2(0, 1)$ -functions g with boundary condition*

$$\begin{aligned} g(0) + g(1) &= 0 \\ g'(0) + g'(1) &= b g(0) \end{aligned} \quad (3.20)$$

for fixed $b \geq 0$. Each b qualifies one of such extensions, with a one-to-one correspondence. Such extensions are ordered with increasing b . The choice $b = 0$ corresponds to anti-periodic boundary conditions.

The application of Theorems 2.4 and 2.5 thus allowed for a fast identification of all non-Friedrichs extensions with Friedrichs lower bound of the minimally defined Laplacian on $[0, 1]$, which would have otherwise required a tedious computation, by means of (3.15)–(3.18), of all the discrete spectra of the various extensions, in order to select those with bottom equal to π^2 .

For completeness, here is how the direct check would have proceeded. Let us limit the analysis to the eigenvalue problem for a generic self-adjoint extension of type (3.16) with the choice (3.19), namely

$$\begin{cases} -g'' = \lambda g \\ g(0) + g(1) = 0 \\ g'(0) + g'(1) = \frac{1}{3}(t - 12) g(0) \\ (\lambda \in \mathbb{R}, g \in H^2(0, 1)) \end{cases} \quad (3.21)$$

for fixed $t \in \mathbb{R}$. g must be of the form $g(x) = A \cos \sqrt{\lambda}x + B \sin \sqrt{\lambda}x$, $A, B \in \mathbb{C}$, and for sure the pairs (g, λ) with

$$g(x) = \sin((2n + 1)\pi x), \quad \lambda = (2n + 1)^2\pi^2, \quad n \in \mathbb{N}_0 \quad (3.22)$$

solve (3.21), showing that all such extensions have the eigenvalues $(2n + 1)^2\pi^2$, $n \in \mathbb{N}_0$, in common. The remaining (i.e., non-sin-only) solutions to (3.21) are obtained

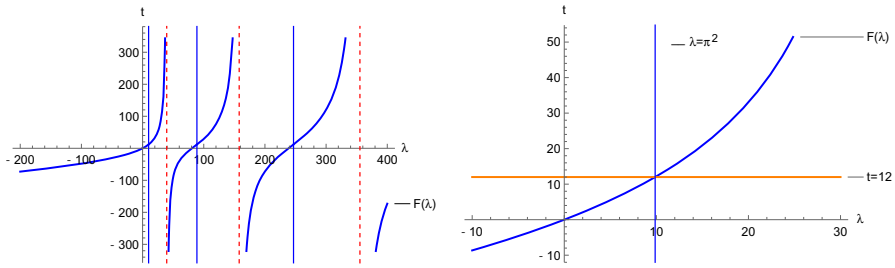


Fig. 1 Left: plot of $F(\lambda)$ defined in (3.23) (blue curves) as compared to the eigenvalues of type (3.22) (vertical blue lines) and the Friedrichs eigenvalues (dashed red lines). Right: magnification of the first positive interval of definition of $F(\lambda)$. Eigenvalues determined by (3.22) and (3.23) correspond to the intersections of the blue curves with the horizontal lines at level t (Color figure online)

imposing $B \neq 0$, and it is then simple to conclude that the admissible λ 's are the (t -dependent) roots of

$$F(\lambda) = t, \quad \text{where} \quad F(\lambda) := 12 - 6\sqrt{\lambda} \frac{1 + \cos \sqrt{\lambda}}{\sin \sqrt{\lambda}} \quad (3.23)$$

(and understanding the above trigonometric functions as hyperbolic functions when $\lambda < 0$). As $F(\lambda)$ increases with λ in all intervals in which it is defined, and $F(\pi^2) = 12$, one deduces that only for $t \geq 12$ the admissible λ 's selected by (3.23) satisfy $\lambda \geq \pi^2$ (see Fig. 1). The spectrum thus determined from (3.22) and (3.23) indeed confirms, by direct inspection, what found in Prop. 3.4 by means of our Theorem 2.5.

3.2 Schrödinger Quantum Particle in \mathbb{R}^3 with Point Interaction

This is an example with deficiency index equal to 1. With respect to the Hilbert space $\mathcal{H} = L^2(\mathbb{R}^3)$ we consider the operator

$$\mathcal{D}(\tilde{S}) = C_0^\infty(\mathbb{R}^3 \setminus \{0\}), \quad \tilde{S} = -\Delta. \quad (3.24)$$

\tilde{S} is densely defined and symmetric, with $m(\tilde{S}) = 0$.

The self-adjoint extensions of \tilde{S} are Hamiltonians for a quantum particle in three dimensions subject to a point interaction supported at $x = 0$, and they are very well studied and understood.

Theorem 3.5 [See, e.g., [1, Chapt. I.1].]

(i) \tilde{S} has unit deficiency index. The Friedrichs extension of \tilde{S} is the self-adjoint (negative) Laplacian on $L^2(0, 1)$ with domain $H^2(0, 1)$. All other self-adjoint extensions of \tilde{S} form the family $\{-\Delta_\alpha \mid \alpha \in \mathbb{R}\}$, where

$$\begin{aligned} \mathcal{D}(-\Delta_\alpha) &= \left\{ g = \phi + \frac{\phi(0)}{\alpha + \frac{1}{4\pi}} G_1 \mid \phi \in H^2(\mathbb{R}^3) \right\} \\ (-\Delta_\alpha + \mathbb{1})g &= (-\Delta + \mathbb{1})\phi \end{aligned} \quad (3.25)$$

and

$$G_1 := (2\pi)^{\frac{3}{2}} \frac{e^{-|x|}}{4\pi|x|}. \quad (3.26)$$

(ii) For each $\alpha \in \mathbb{R}$,

$$\sigma_{\text{ess}}(-\Delta_\alpha) = \sigma_{\text{ac}}(-\Delta_\alpha) = [0, +\infty), \quad \sigma_{\text{sc}}(-\Delta_\alpha) = \emptyset, \quad (3.27)$$

and

$$\sigma_p(-\Delta_\alpha) = \begin{cases} \emptyset & \text{if } \alpha \in [0, +\infty) \\ \{-(4\pi\alpha)^2\} & \text{if } \alpha \in (-\infty, 0). \end{cases} \quad (3.28)$$

The negative eigenvalue $-(4\pi\alpha)^2$, when it exists, is simple and the corresponding eigenfunction is $|x|^{-1}e^{4\pi\alpha|x|}$.

We see from Theorem 3.5 that \tilde{S} admits a collection of non-Friedrichs extensions with Friedrichs lower bound, and precisely

$$m(-\Delta_\alpha) = 0 = m(\tilde{S}) \quad \forall \alpha \geq 0. \quad (3.29)$$

In order to recover such a conclusion from the abstract setting of Sect. 2, let us consider

$$S := \tilde{S} + \mathbb{1}. \quad (3.30)$$

Clearly, $m(S) = 1$. The self-adjoint extensions of \tilde{S} and of S then only differ by a trivial shift. As we intend to analyse the extensions of S within the extension scheme of Theorem 2.1, rather than using von Neumann's extension theorem as in [1], let us follow closely the discussion made in [10, Sect. 3], where indeed the Kreĭn-Višik-Birman scheme was employed.

We shall denote, respectively, by $\hat{\cdot}$ and \cdot the Fourier and inverse Fourier transform $L^2(\mathbb{R}^3, dx) \rightarrow L^2(\mathbb{R}^3, dp)$ with the convention

$$\hat{f}(p) = \frac{1}{(2\pi)^{\frac{3}{2}}} \int_{\mathbb{R}^3} e^{-ip \cdot x} f(x) dx.$$

In particular,

$$G_1 = (2\pi)^{\frac{3}{2}} \frac{e^{-|x|}}{4\pi|x|} = \left(\frac{1}{p^2 + 1} \right)^\vee. \quad (3.31)$$

It is possible to prove the following.

Theorem 3.6 [10, Sect. 3](i) S has deficiency space

$$\ker S^* = \text{span}\{G_1\}. \quad (3.32)$$

(ii) The Friedrichs extension of S is the operator

$$\begin{aligned} \mathcal{D}(S_F) &= \left\{ g \in L^2(\mathbb{R}^3) \mid \begin{array}{l} \widehat{g} = \widehat{f} + (p^2 + 1)^{-1}\eta \\ f \in \mathcal{D}(\overline{S}), \eta \in \mathbb{C} \end{array} \right\} \\ \widehat{S_F g} &= (p^2 + 1)\widehat{g}. \end{aligned} \quad (3.33)$$

(iii) All other self-adjoint extensions of S are of the form S_t for some $t \in \mathbb{R}$, where

$$\begin{aligned} \mathcal{D}(S_t) &= \left\{ g \in L^2(\mathbb{R}^3) \mid \begin{array}{l} \widehat{g} = \widehat{f} + (p^2 + 1)^{-2}t\xi + (p^2 + 1)^{-1}\xi \\ f \in \mathcal{D}(\overline{S}), \xi \in \mathbb{C} \end{array} \right\} \\ \widehat{S_t g} &= (p^2 + 1)(\widehat{f} + (p^2 + 1)^{-2}t\xi). \end{aligned} \quad (3.34)$$

This is precisely formula (2.1) of Theorem 2.1 specialised to the case where $\ker S^*$ is one-dimensional and T is therefore the operator of multiplication by the real number t .

(iv) One has

$$S_t = -\Delta_\alpha + \mathbb{1} \quad \text{for} \quad \alpha = \frac{t-2}{8\pi}. \quad (3.35)$$

Clearly, $S_F - m(S)\mathbb{1} = S_F - \mathbb{1} = \widetilde{S}_F$, the self-adjoint (negative) Laplacian on $L^2(\mathbb{R}^3)$. Therefore, unlike the example discussed in Sect. 3.1,

$$\ker(S_F - \mathbb{1}) = \{0\}. \quad (3.36)$$

$S_F - \mathbb{1}$ is then invertible on its range and so are the powers $(S_F - \mathbb{1})^\delta$, $\delta > 0$. On such a space, $(S_F - \mathbb{1})^{-\delta}$ acts, in Fourier transform, as the multiplication by $|p|^{-2\delta}$.

The analogue of Lemma 3.1 is now the following.

Lemma 3.7 *One has*

$$V := \text{ran}(S_F - \mathbb{1})^{\frac{1}{2}} \cap \ker S^* = \text{span}\{G_1\} \quad (3.37)$$

and

$$(S_F - \mathbb{1})^{-\frac{1}{2}} G_1 = \left(\frac{1}{|p|(p^2 + 1)} \right)^\vee. \quad (3.38)$$

Proof The fact that $G_1 \in \ker S^*$ is stated in Theorem 3.6(i). As

$$\frac{1}{|p|(p^2 + 1)} \in L^2(\mathbb{R}^3, dp)$$

and

$$(S_F - \mathbb{1})^{\frac{1}{2}} \left(\frac{1}{|p|(p^2 + 1)} \right)^\vee = \left(\frac{1}{(p^2 + 1)} \right)^\vee = G_1,$$

hence $G_1 \in \text{ran}(S_F - \mathbb{1})^{\frac{1}{2}}$. V can be at most one-dimensional, thus (3.37) is proved, and so is (3.38) as well. \square

Owing to Lemma 3.7, Theorem 2.4 is applicable: S admits non-Friedrichs extensions with Friedrichs lower bound, and so does therefore \tilde{S} , consistently with what previously observed in (3.29).

Furthermore, with the notation of Theorem 2.5,

$$\mathcal{D}[q] = \mathcal{D}(T_q) = V = \text{span}\{G_1\}, \quad (3.39)$$

thus T_q is an operator of multiplication by some real number t_q ,

$$T_q G_1 = t_q G_1. \quad (3.40)$$

Since

$$\begin{aligned} \langle G_1, T_q G_1 \rangle &= \|G_1\|_2^2 + \|(S_F - \mathbb{1})^{-\frac{1}{2}} G_1\|_2^2 \\ &= \left\| \frac{1}{p^2 + 1} \right\|_2^2 + \left\| \frac{1}{|p|(p^2 + 1)} \right\|_2^2 \\ &= \pi^2 + \pi^2 \\ &= 2 \|G_1\|_2^2 \end{aligned}$$

(having used (3.38) in the second identity), then necessarily $t_q = 2$.

Theorem 2.5, in parts (iii) and (iv), then states that the self-adjoint extensions S_T of S with $\mathfrak{m}(S_T) = \mathfrak{m}(S_F)$ are those labelled by self-adjoint operators T with $T \geq T_q$. Such T 's, apart from the one parametrising the Friedrichs extension, are therefore such that

$$\begin{aligned} \mathcal{D}(T) &= V = \text{span}\{G_1\} \\ T &\text{ is the multiplication by some } t \geq 2. \end{aligned} \quad (3.41)$$

For what argued in Theorem 3.6(iii), such extensions are precisely the operators S_t that one reads out from formula (3.34) with $t \geq 2$. In turn, the correspondence formula (3.35) leads to the conclusion that the self-adjoint extensions of \tilde{S} with Friedrichs lower bound are precisely those $-\Delta_\alpha$'s with $\alpha \geq 0$.

3.3 Radial Problem in Hydrogenoid-Like Hamiltonians

It is worth mentioning another example with unit deficiency index, even without working out here the steps through which Theorems 2.4 and 2.5 are applied, which are in fact completely analogous to the computations of Sects. 3.1 and 3.2.

For given $\nu \in \mathbb{R}$, let us now consider

$$\mathcal{D}(S_\nu) = C_0^\infty(\mathbb{R}^+), \quad S_\nu = -\frac{d^2}{dx^2} + \frac{\nu}{x}, \quad (3.42)$$

a densely defined and symmetric operator on the Hilbert space $\mathcal{H} = L^2(\mathbb{R}^+)$ with lower bound $m(S_\nu) = 0$. One typical emergence of S_ν in mathematical physics is as the minimally defined zero-momentum radial operator in the construction of a quantum hydrogenoid Hamiltonian with an additional point interaction at the center of the Coulomb potential: S_ν is indeed well known and thoroughly studied, and we refer to [4, Sect. 1.4] and references therein for an updated historical overview.

Hardy's inequality implies that S_ν is lower semi-bounded, and in particular obviously

$$m(S_\nu) = 0 \quad \forall \nu \geq 0 \quad (3.43)$$

(repulsive Coulomb interaction). A standard limit-point limit-circle argument shows that S_ν has unit deficiency index. Its self-adjoint extensions are studied in the literature by means of various extension schemes, including recently in [4] by means of the general Theorem 2.1 above.

Theorem 3.8 [4, Theorems 2 and 4].

(i) *The self-adjoint extensions of S_ν in $L^2(\mathbb{R}^+)$ form the family $\{S_\nu^{(\alpha)} \mid \alpha \in \mathbb{R} \cup \{\infty\}\}$, where $\alpha = \infty$ labels the Friedrichs extension, and*

$$\begin{aligned} \mathcal{D}(S_\nu^{(\alpha)}) &= \left\{ g \in L^2(\mathbb{R}^+) \mid \begin{array}{l} -g'' + \frac{\nu}{r}g \in L^2(\mathbb{R}^+) \\ \text{and } g_1 = 4\pi\alpha g_0 \end{array} \right\} \\ S_\nu^{(\alpha)} g &= -g'' + \frac{\nu}{r}g, \end{aligned} \quad (3.44)$$

g_0 and g_1 being the existing limits

$$\begin{aligned} g_0 &:= \lim_{r \downarrow 0} g(r) \\ g_1 &:= \lim_{r \downarrow 0} r^{-1}(g(r) - g_0(1 + \nu r \ln r)). \end{aligned} \quad (3.45)$$

(ii) *For $\nu > 0$ one has*

$$\sigma_{ess}(S_\nu^{(\alpha)}) = \sigma_{ac}(S_\nu^{(\alpha)}) = [0, +\infty), \quad \sigma_{sc}(S_\nu^{(\alpha)}) = \emptyset, \quad (3.46)$$

and

$$\sigma_p(S_v^{(\alpha)}) = \begin{cases} \emptyset & \text{if } \alpha \geq \alpha_v \\ \{E_+^{(v,\alpha)}\} & \text{if } \alpha < \alpha_v, \end{cases} \tag{3.47}$$

where

$$\alpha_v := \frac{\nu}{4\pi} (\ln \nu + 2\gamma - 1) \tag{3.48}$$

($\gamma \sim 0.577$ being the Euler-Mascheroni constant) and $E_+^{(v,\alpha)}$ is the only simple negative root of $\mathfrak{F}_\nu(E) = \alpha$ with

$$\mathfrak{F}_\nu(E) := \frac{\nu}{4\pi} \left(\psi \left(1 + \frac{\nu}{2\sqrt{|E|}} \right) + \ln(2\sqrt{|E|}) + 2\gamma - 1 - \frac{\sqrt{|E|}}{\nu} \right) \tag{3.49}$$

($\psi(z) = \Gamma'(z) / \Gamma(z)$ being the digamma function).

When $\nu > 0$ Theorem 3.8 thus shows that

$$\begin{aligned} m(S_v^{(\alpha)}) &= E_+^{(v,\alpha)} < 0 = m(S_\nu) & \text{if } \alpha < \alpha_v \\ m(S_v^{(\alpha)}) &= 0 = m(S_\nu) & \text{if } \alpha \geq \alpha_v, \end{aligned} \tag{3.50}$$

yet another example of the presence of a sub-class of non-Friedrichs extensions with Friedrichs lower bound.

Also on this example it is easy to test the applicability of our Theorems 2.4-2.5. As done in Sect. 3.2, since $m(S_\nu) = 0$, a positive shift must be performed first. For similar purposes the analysis of the shifted operator

$$S_\nu := S_\nu + \frac{\nu^2}{4\kappa^2} \mathbb{1} \quad (\kappa \in \mathbb{R}) \tag{3.51}$$

and of its self-adjoint extensions was worked in [4, Sect. 2], which we refer to for the details. The special value of the shift (3.51) was chosen in [4] in order to be able to solve the ODE $S_\nu^* u = 0$ by means of special functions, this way characterising explicitly the deficiency space $\ker S_\nu^*$. The Friedrichs extension $S_{\nu,F}$ of S_ν was also characterised in [4, Sect. 2]. This provides all the ingredients to investigate the intersection (2.9) and apply Theorems 2.4 and 2.5 so as to reproduce (3.50).

3.4 Schrödinger Quantum Particle on an Interval, Continued

In this last example let us show how our abstract results apply also to the edge case in which no other self-adjoint extension has the same bottom as the Friedrichs'.

Acting on the Hilbert space $\mathcal{H} = L^2(0, 1)$, let us consider the densely defined and symmetric operator

$$S = -\frac{d^2}{dx^2}, \quad \mathcal{D}(S) = \{f \in H^2(0, 1) \mid f(0) = f'(0) = f(1) = 0\}. \quad (3.52)$$

Again by Poincaré inequality, S is semi-bounded from below with $m(S) = \pi^2$.

It is standard to see that S is a symmetric extension of the operator considered in Sect. 3.1, but now with unit deficiency index, with deficiency subspace

$$\ker S^* = \overline{\text{span}\{\mathbf{1} - x\}}, \quad (3.53)$$

and with Friedrichs extension S_F given exactly by the already considered operator (1.5). Thus, (3.2)–(3.3) above apply, and therefore the spanning function $\mathbf{1} - x$ cannot belong to $\text{ran}(S_F - \pi^2 \mathbb{1})^{\frac{1}{2}} = \text{ran}(S_F - \pi^2 \mathbb{1}) = \{\sin \pi x\}^\perp$, for

$$\int_0^1 (1-x) \sin \pi x \, dx = \frac{1}{\pi} \neq 0. \quad (3.54)$$

The conclusion is then

$$\text{ran}(S_F - \pi^2 \mathbb{1})^{\frac{1}{2}} \cap \ker S^* = \{0\}. \quad (3.55)$$

Theorem 2.4 then predicts that S_F is the only self-adjoint extension of S with lower bound π^2 .

For an independent direct check of the latter conclusion, one can argue that the one-parameter family of self-adjoint extensions of S consists of negative Laplacians on H^2 -functions g subject to the condition (3.16) with $c = 0$ and $b_1 \in \mathbb{R}$, i.e.,

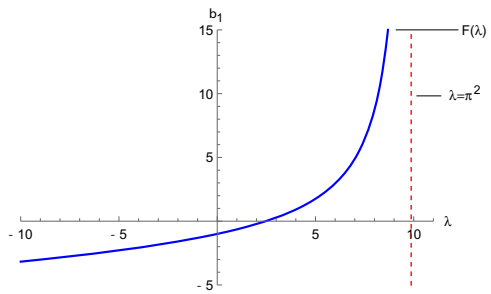
$$g'(0) = b_1 g(0), \quad g(1) = 0. \quad (3.56)$$

The corresponding eigenvalue problem, namely the analogue of (3.21) now with the boundary condition (3.56), is solved by the roots $\lambda \in \mathbb{R}$ of

$$-\frac{\sqrt{\lambda}}{\tan \sqrt{\lambda}} = b_1 \quad (3.57)$$

for given $b_1 \in \mathbb{R}$ that qualifies each extension (and understanding (3.57) in terms of the hyperbolic tangent when $\lambda < 0$). One easily sees (Fig. 2) that (3.57) always admits a solution $\lambda < \pi^2 = m(S_F)$, and only asymptotically for $b_1 \rightarrow +\infty$ does the smallest root approach π^2 , thus proving that each such extension is strictly lower than S_F .

Fig. 2 Graphic solution to (3.57)



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